

# Neutrinos and the large scale structure

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- the role of neutrinos in the cosmic budget
- neutrinos and the formation of cosmic structures
- cosmological constraints on the neutrino mass

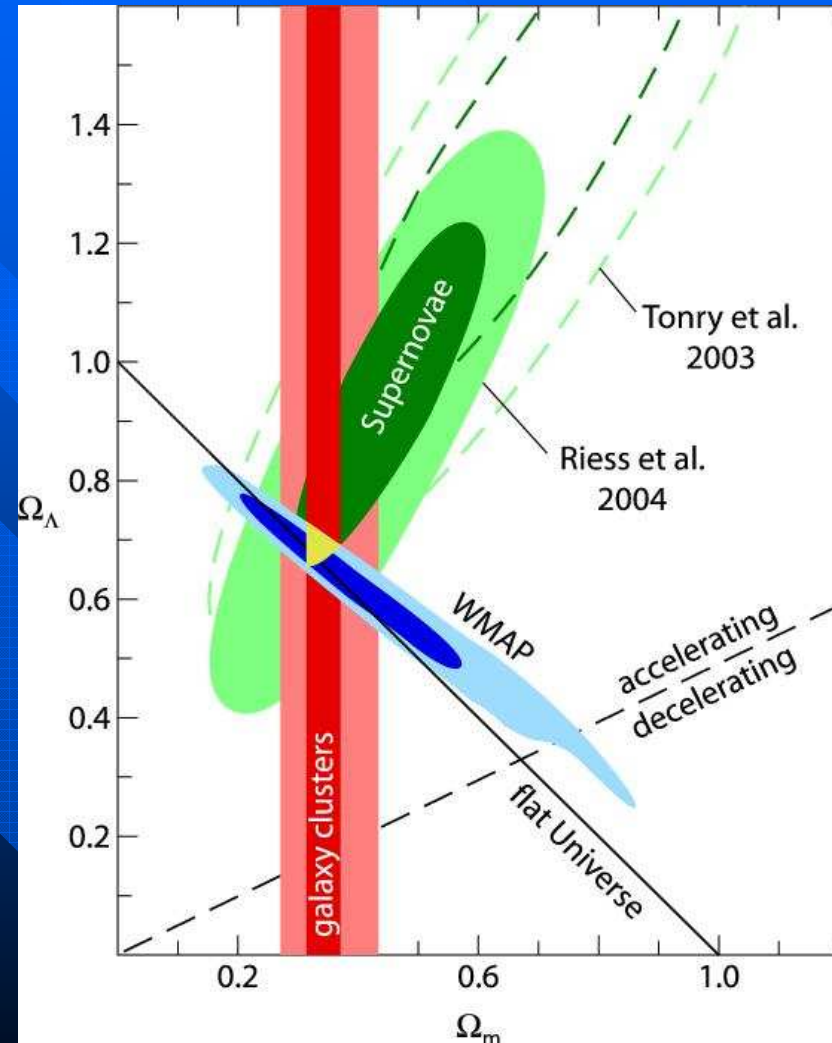
**Perspectives in Neutrino Physics & Astrophysics**  
**Bologna, 17th June 2005**

# The cosmic budget

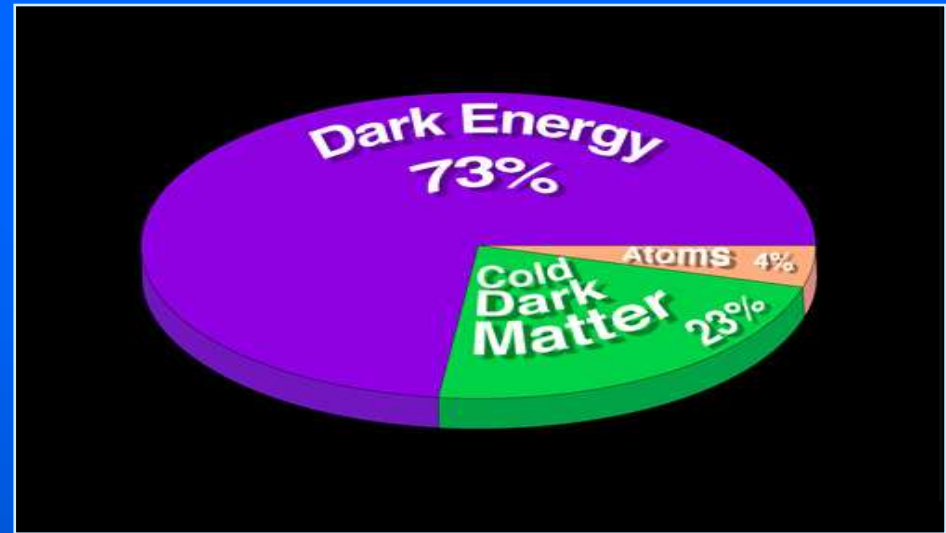
Independent data sets give a consistent determination of the amount of Dark Energy and Dark Matter in the Universe. The relative weights being measured by their density parameter

$$\Omega_i = \rho_i / \rho_c$$

where  $\rho_c = \cong 10^{-29} \text{ g/cm}^3$  is the *critical density* i.e. the energy density which closes the Universe



# The cosmic budget



- ✓ About 73% of the energy content of our Universe is in the form of some exotic component, called **Dark Energy**, or "**Quintessence**", which causes a large-scale cosmic repulsion among celestial objects, thereby mimicking a sort of anti-gravity effect. The simplest dark energy candidate is the Cosmological Constant  $\Lambda$ .
- ✓ Only about 4% of the cosmic energy budget is in the form of ordinary "baryonic" matter, out of which only a small fraction shines in the galaxies (quite likely most of the baryon reside in filaments forming the **Warm-Hot Intergalactic Medium (WHIM)**, a sort of cosmic web connecting the galaxies and clusters of galaxies).
- ✓ About 23% of the cosmic budget is made of **Dark Matter**, a collisionless component whose presence we only perceive gravitationally.

# Cosmological neutrinos

Neutrinos are in equilibrium with the primeval plasma through weak interaction reactions.

They **decouple** from the plasma at

$$T_{\text{dec}} \approx 1 \text{ MeV}$$

Today we have a cosmological neutrino background at a **temperature**

$$T_{\nu} = (4/11)^{1/3} T_{\gamma} \approx 1.945 \text{ K},$$

corresponding to  $kT_{\nu} \approx 1.68 \cdot 10^{-4} \text{ eV}$

# The Neutrino density

This corresponds to a present neutrino number density of

$$n_{0\nu} \approx 0.1827 \cdot T_\nu^3 \approx 112 \text{ cm}^{-3}$$

That for a massive neutrino translates in

$$\rho_{0\nu} \approx 1.9 \cdot N_\nu \langle m_\nu / 10 \text{ eV} \rangle \cdot 10^{-30} \text{ g cm}^3$$

or equivalently

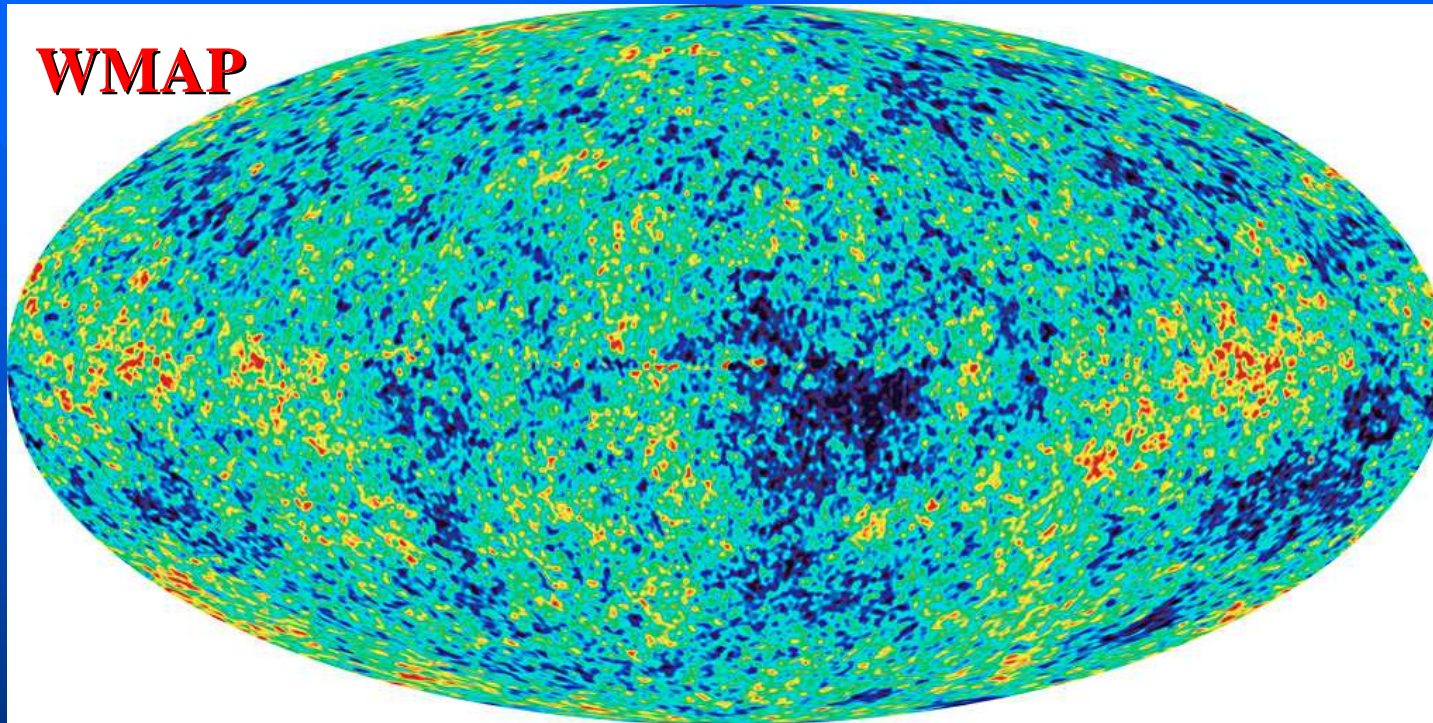
$$\Omega_{0\nu} h^2 \approx 0.1 \cdot N_\nu \langle m_\nu / 10 \text{ eV} \rangle$$

i.e. in order to be a good candidate for the dark matter component of our universe ( $\Omega_{0M} h^2 \approx 0.14$ ), neutrinos need to have a **mean mass of approximately 5 eV!**

A **direct detection is very difficult** but...

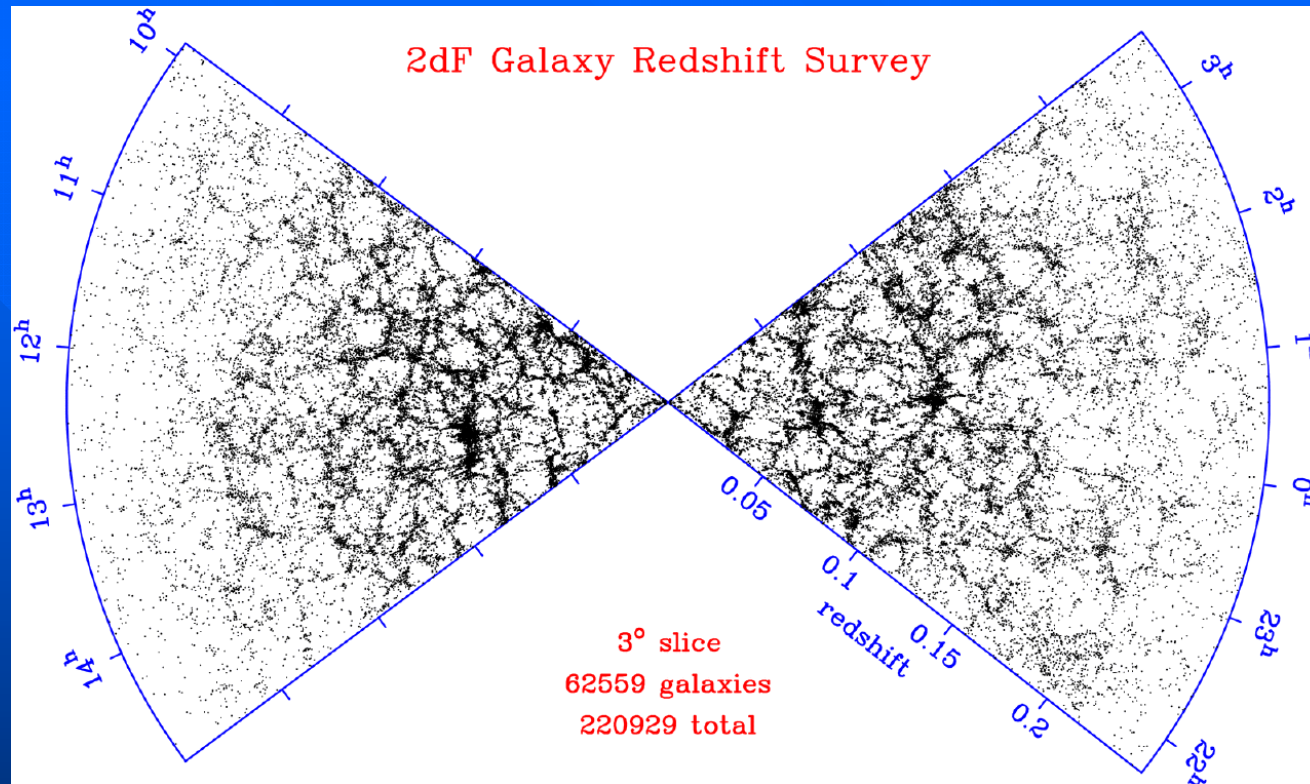
They have a strong impact on the **formation and evolution of cosmic structures**, the so-called **cosmic clustering**, which now can be accurately measured.

# The Model for Structure formation



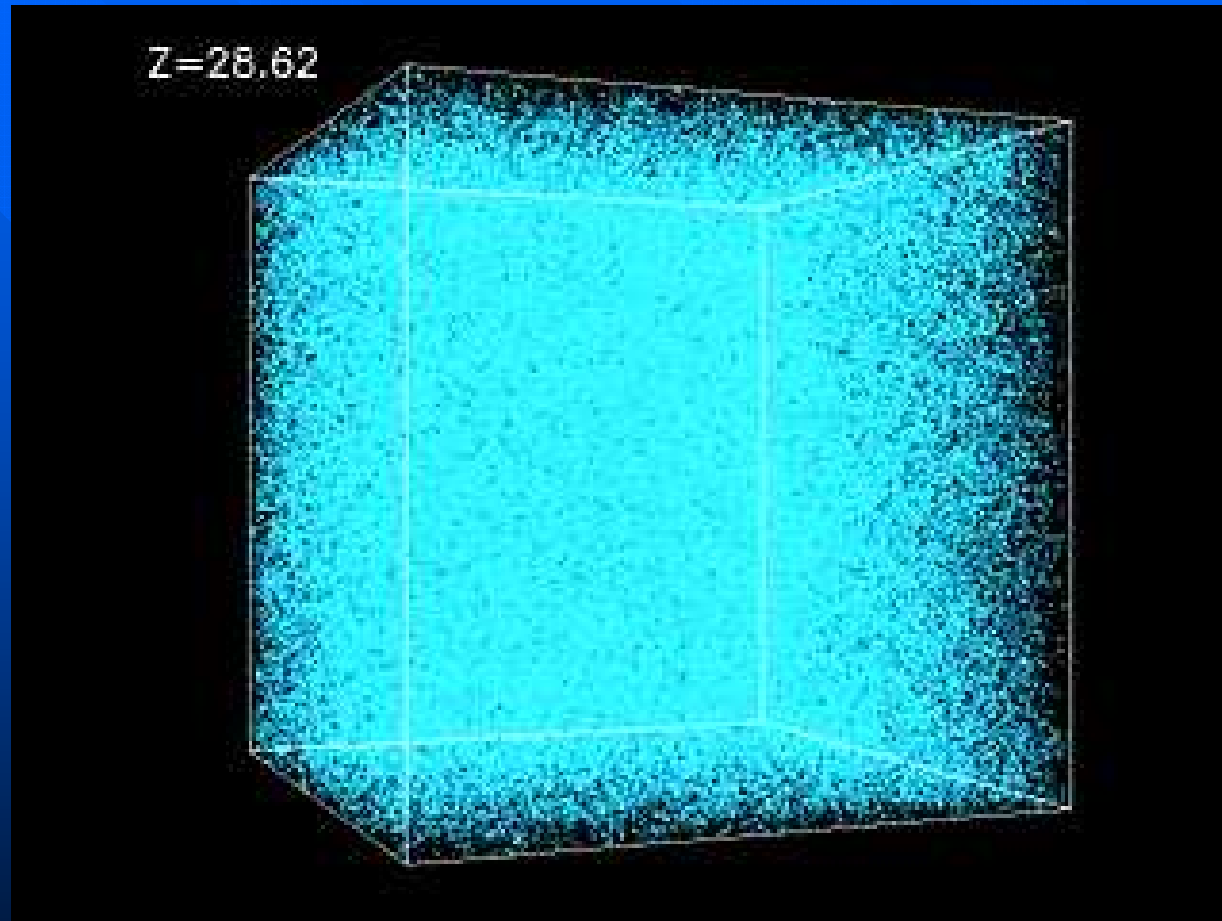
The **cosmic microwave background** (CMB) tells us that the universe is almost perfectly uniform spatially, with density variations from place to place only at the level of  $10^{-5}$ .

# The Model for Structure formation



**Gravitational instability** caused these tiny fluctuations to grow in amplitude into the large scale structure we observe: gravity is an attractive force and tends to increase the overdensity over time

# Redshift evolution of clustering





# The linear solution

**IF** all the matter contributing to the cosmic density is able to cluster (like dark matter or ordinary matter with negligible pressure), then density fluctuations  $\delta$  grow as the cosmic expansion factor  $a \propto (1+z)^{-1}$ , i.e.

$$\delta \propto a,$$

But, **IF** some fraction  $(1-\Omega_*)$  is unable to cluster (i.e. it is gravitationally inert), then the growth will be **slower**

$$\delta \propto a^p,$$

where **p**  $\approx \Omega_*^{0.6}$ .

Note that the inert component can include dark energy if present and photons and neutrinos on sufficiently large scales.

# Consequences

- **At early times**, the cosmic density is dominated by photons. This implies  $p \approx 0$ : fluctuations cannot start growing until the epoch of matter-domination (MD), starting at  $z \approx 3700$
- **At recent times** ( $z \approx 0.3$ ), the density is dominated by dark energy (DE), which gradually stops the growth of fluctuations after a net growth factor of about  $(a_{\text{DE}}/a_{\text{MD}}) \approx 4700$

# What about neutrinos?

Massive non-relativistic neutrinos cannot cluster on small scales because of their high velocities. In the period between matter and dark energy domination, neutrinos are a roughly constant fraction  $f_\nu = (1 - \Omega_*)$  of the matter density.

Then the net fluctuation growth factor is

$$(a_{\text{DE}}/a_{\text{MD}})^p \approx 4700^p \approx 4700 \exp(-4 f_\nu)$$

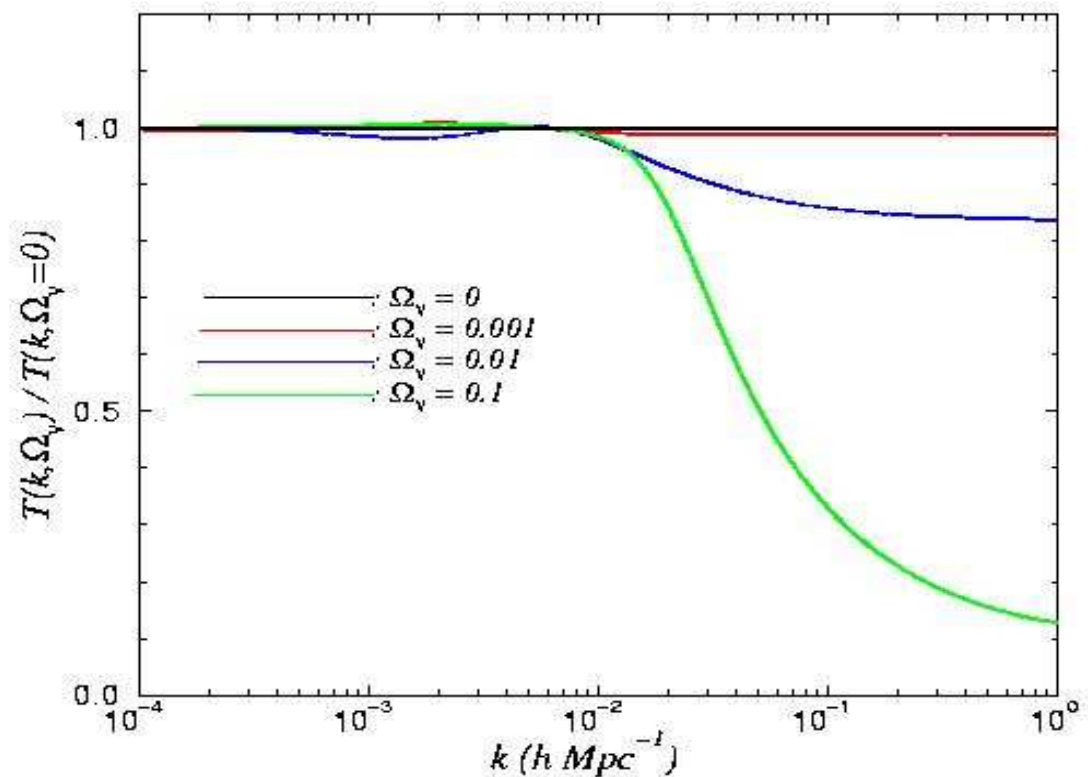
**Even a small neutrino fraction has a large effect!**

# The transfer function $T(k)$

In cosmology this effect can be quantified by using the **density power spectrum  $P(k)$** , giving the variance of fluctuations  $\delta$  in Fourier space. Usually this can be written as

$$P(k) = A k^n T^2(k)$$

**Neutrino free streaming:**  
 $\Delta P(k)/P(k) = -8f_\nu$



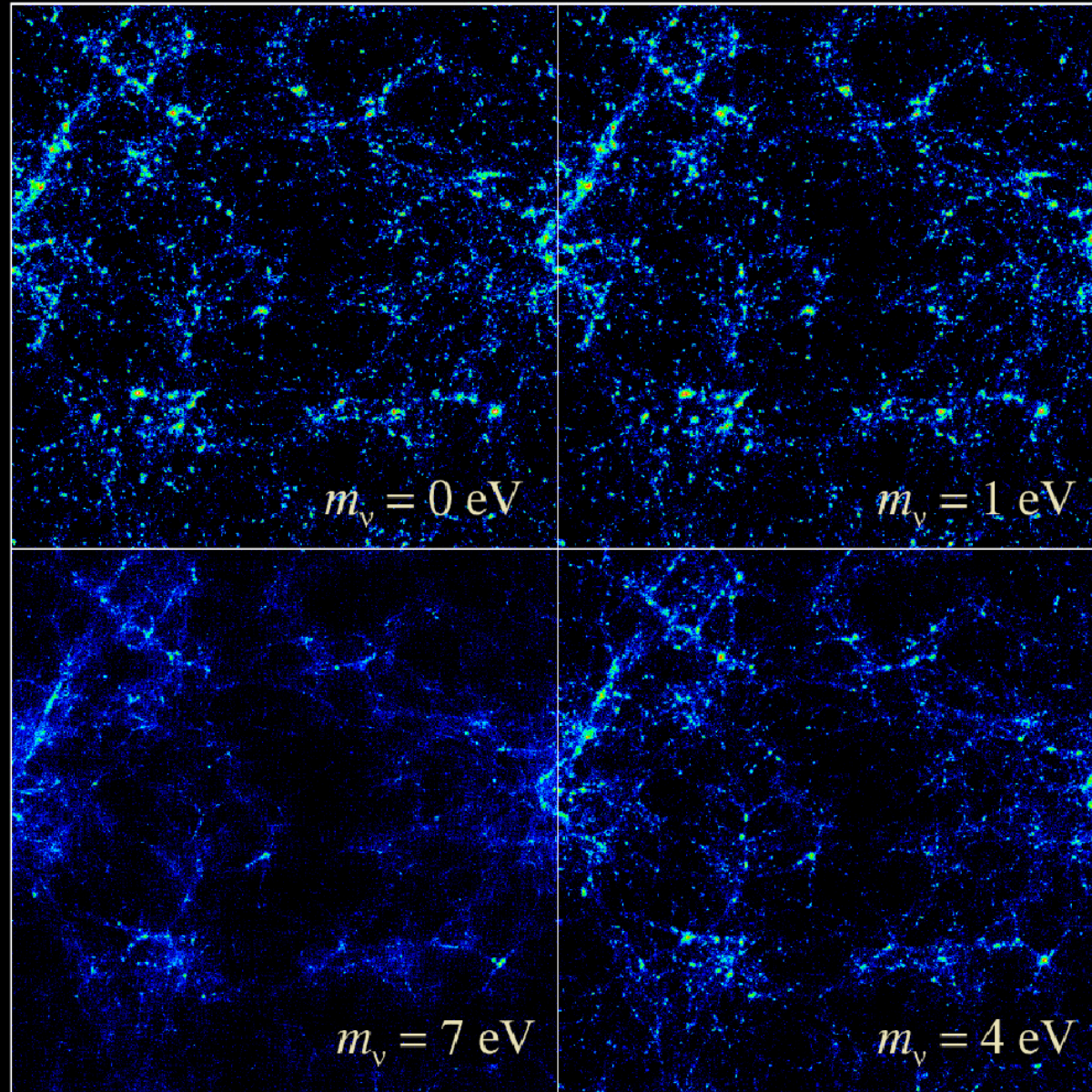
# Practical consequences

- There is a scale, called **neutrino free-streaming scale**, below which clustering is strongly suppressed.
- Neutrinos will not cluster in overdense clumps so small that their escape velocity is much smaller than the typical neutrino velocity.
- **On larger scales** neutrinos behave just as cold dark matter:  $\Omega_* = 1$  and  $p=1$

**The power spectrum changes its shape in a characteristic way**

# N-body results

There is  
less  
clustering  
in models  
with  
massive  
neutrinos



Ma '96

# The top-down scenario

Now we know that  $\Omega_{0m} \approx 0.3$ .

If we assume that all dark matter is contributed by neutrinos, because of free-streaming there will be a strong suppression of power at small scales.

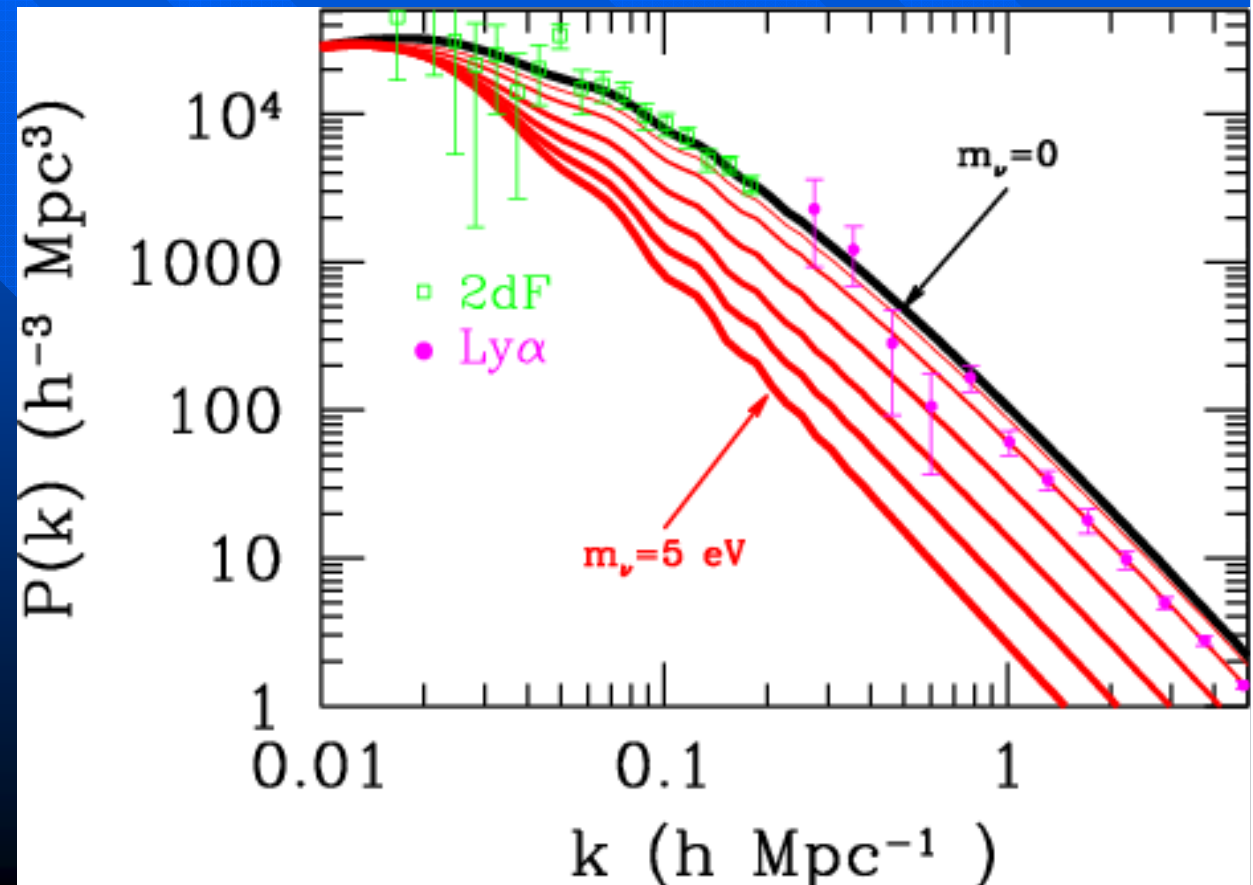
Consequently cosmic structures would have formed first at large scales (galaxy clusters), and smaller structures (like galaxies) would form later by fragmentation:

this is the so-called **top-down scenario**

But, starting from late '80s, we have evidences in favour of a **bottom-up** structure formation (hierarchical) model, where objects formed first at small scale.

Now this is confirmed by observational data. A **cold** (i.e. non-relativistic when it decoupled from the thermal background) **dark matter** (CDM) component is strongly favoured

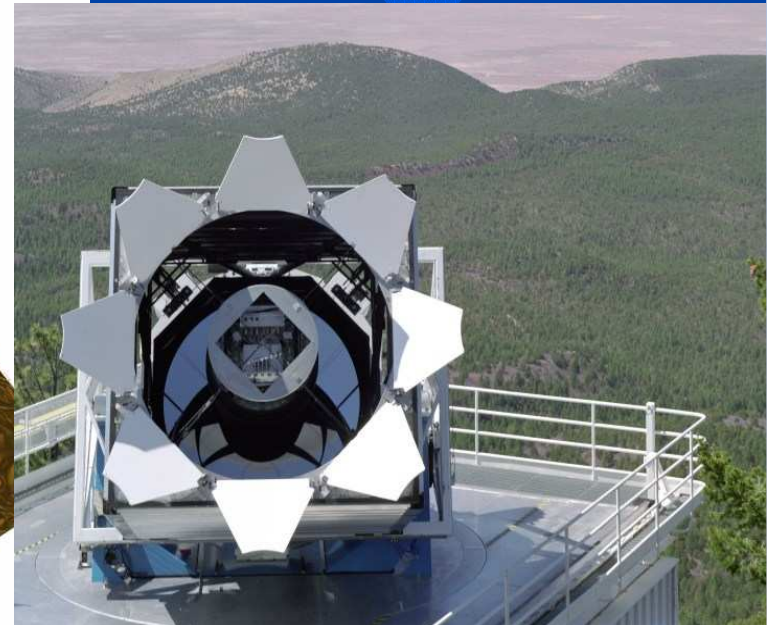
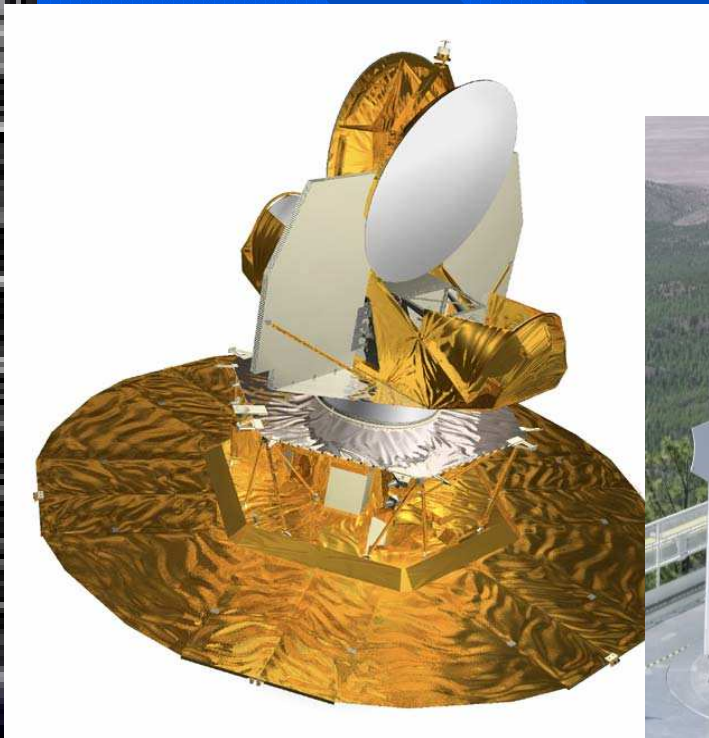
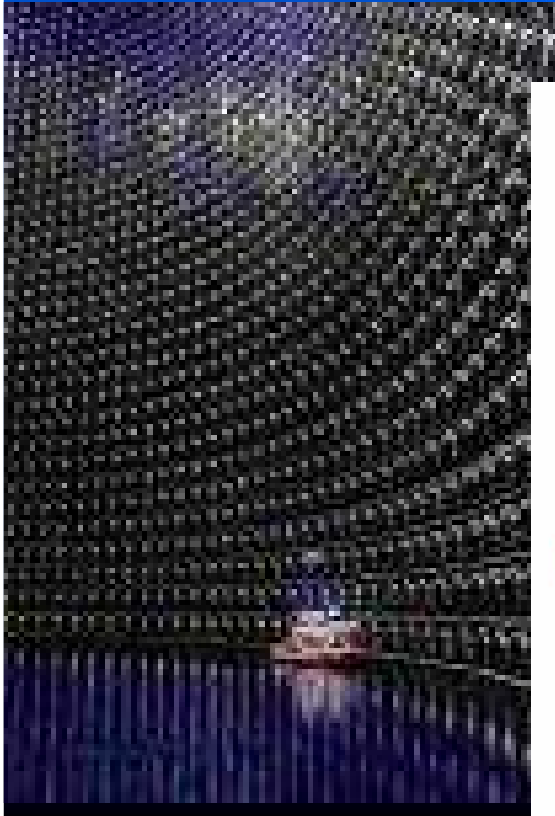
**dark matter cannot be dominated by neutrinos!**





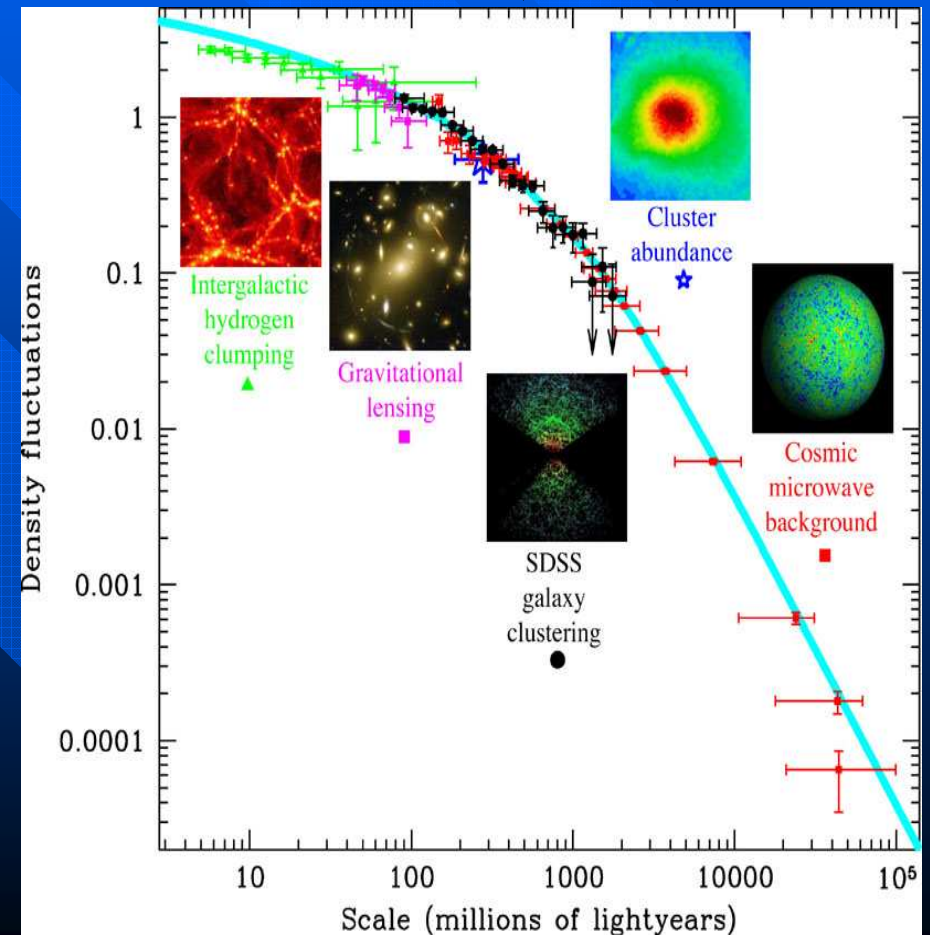
# Weighing neutrinos

However, neutrinos, even if non-dominant, are **massive and abundant**. So, if we have accurate measurements of cosmic clustering (as we start to have now), we can hope to use cosmological observations to put constraints on the neutrino mass which can be combined with laboratory bounds



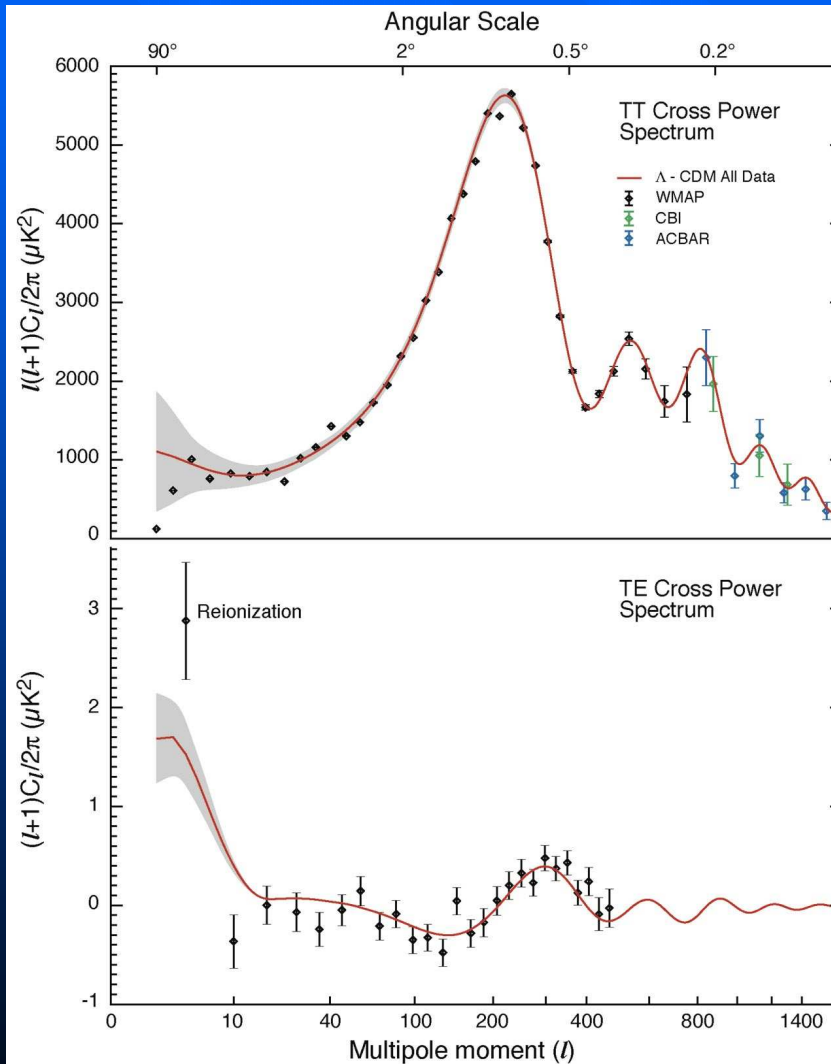
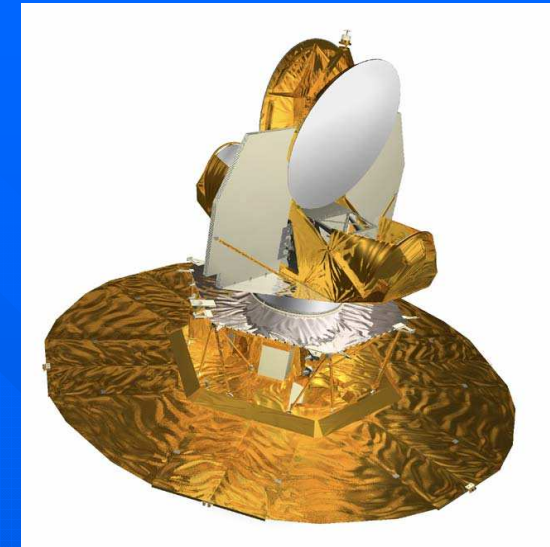
# Cosmological observables

- Cosmic microwave background (CMB)
- Galaxy surveys & large scale structure (LSS)
- Lyman alpha forest
- Galaxy clusters
- Gravitational lensing
- ...



Tegmark

# WMAP CMB anisotropies

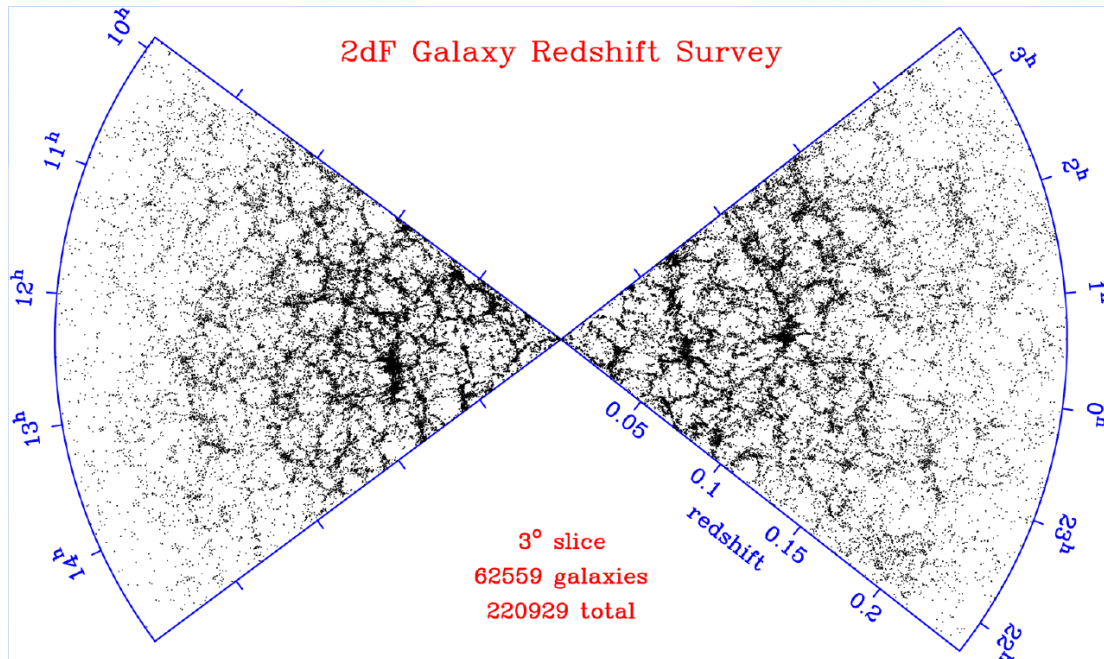


**CMB alone is NOT sensitive to massive neutrinos:** there is only a small enhancement of the acoustic peaks.

However, they are able to put strong constraints on the matter density and on other parameters: this allows, when combined with other data, to

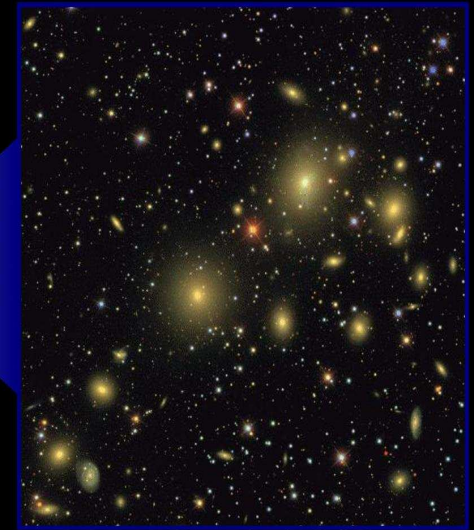
**break degeneracies**

# Galaxy surveys

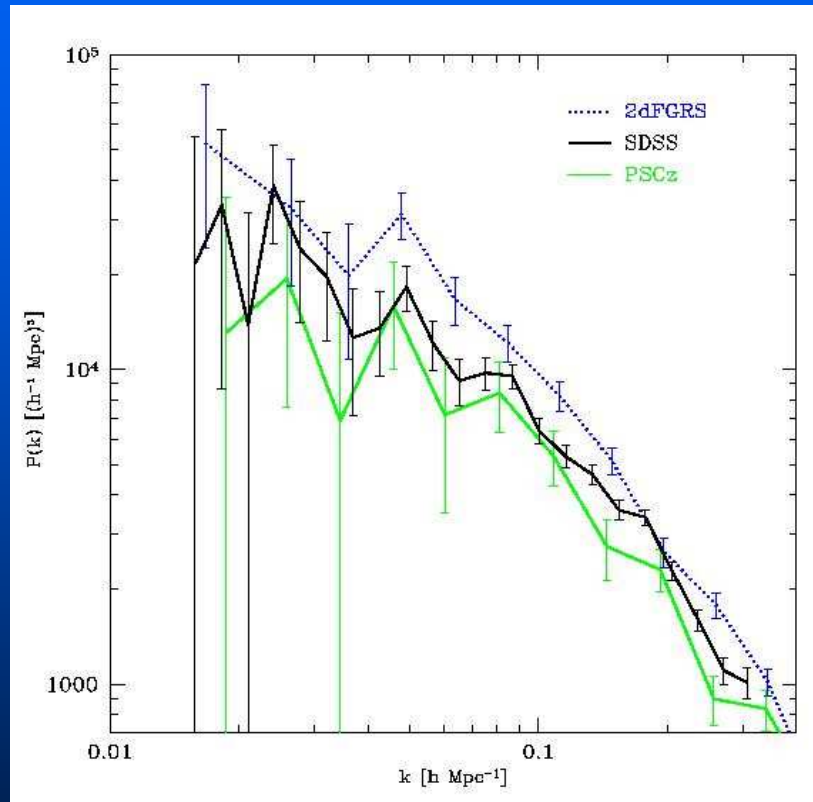


- Large surveys with >200k galaxy redshifts: **2dF** and **SDSS**
- In linear regime, sensitive to neutrino fraction  $f_\nu = \Omega_\nu / \Omega_m$

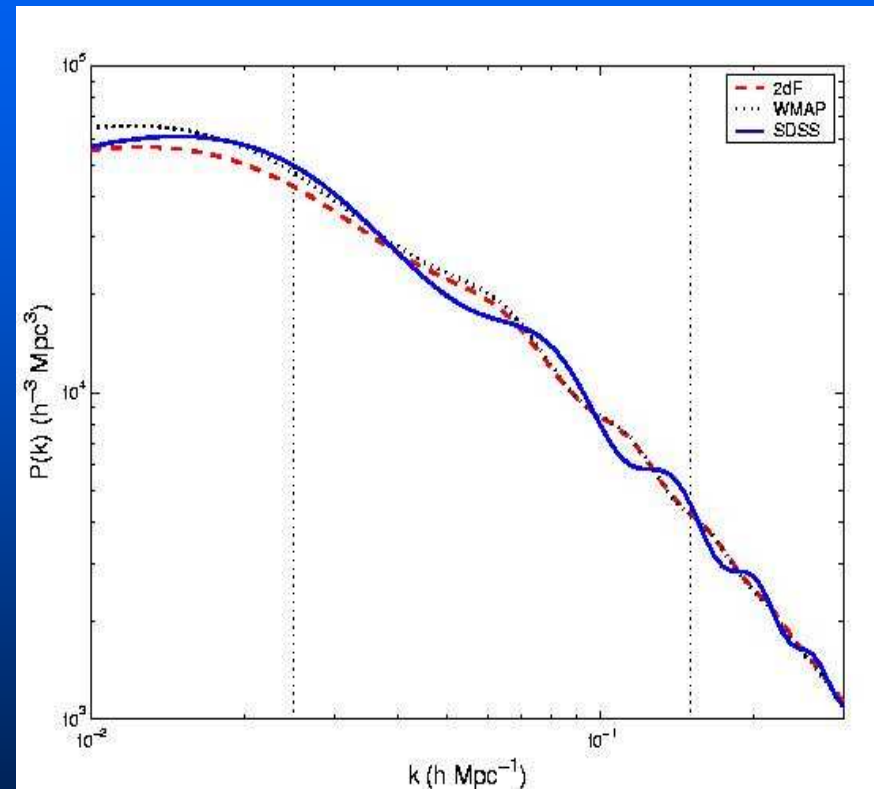
**SDSS**



# 2dF vs SDSS Power spectra

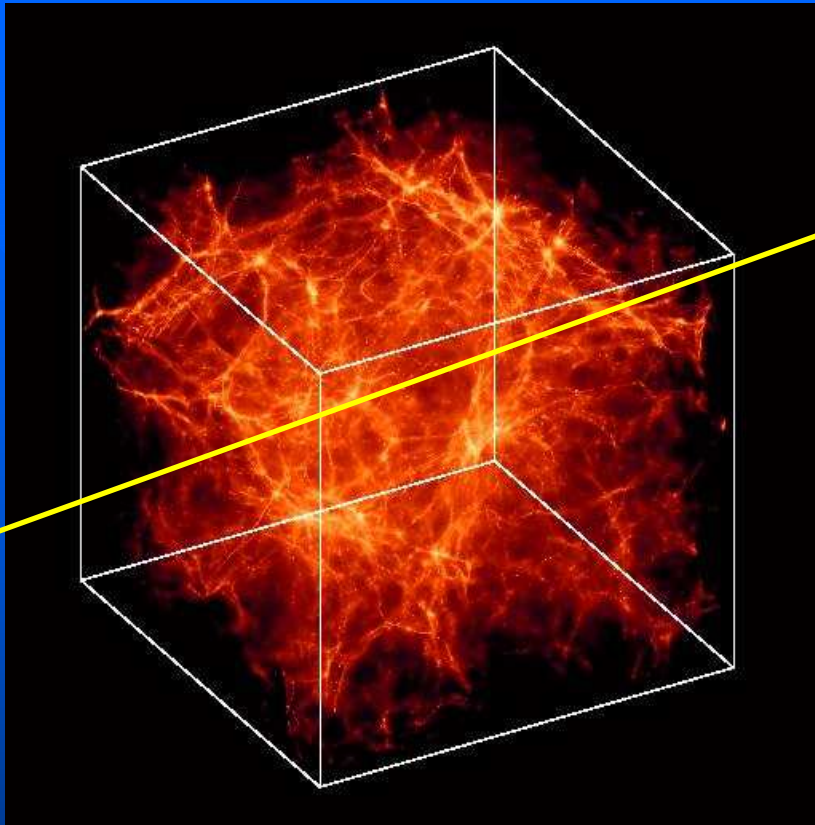
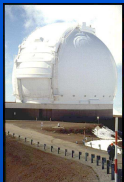


Tegmark et al. 2003

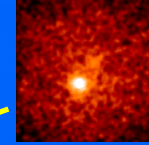


Pope et al. 2004

obs.



QSO



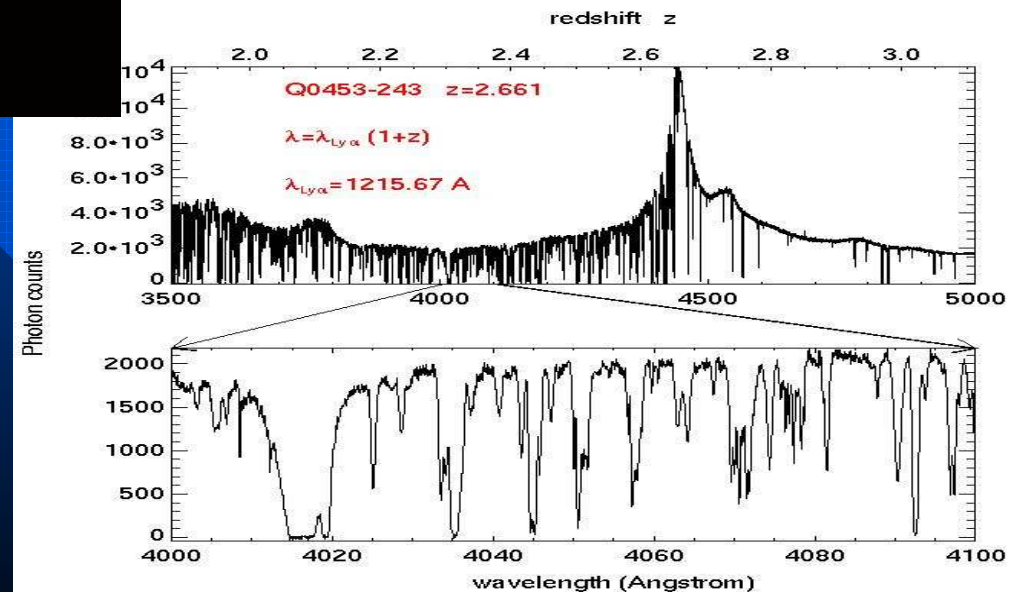
Lyman- $\alpha$   
forest

80 % of the baryons at  $z=3$   
are in the Lyman- $\alpha$  forest  
(Rauch 1998)

baryons as tracer of the dark  
matter density field

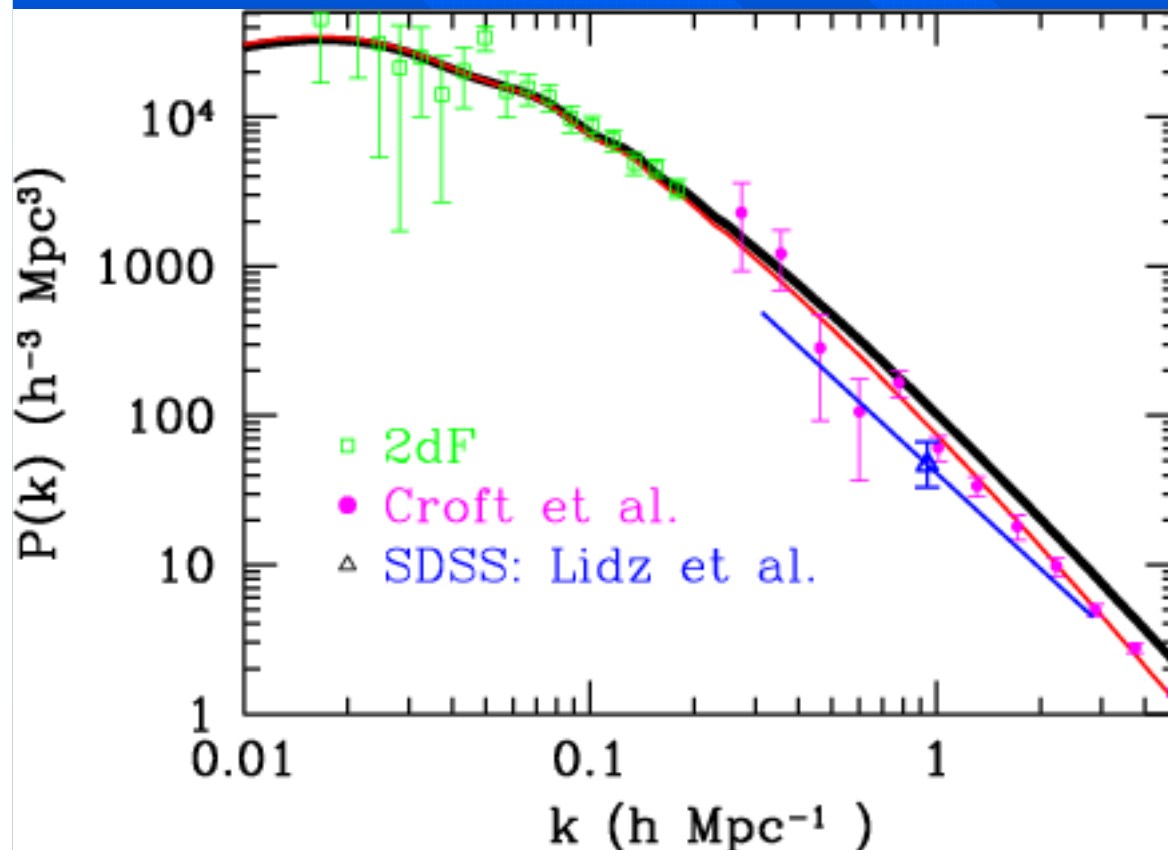
$\delta_{\text{IGM}} \sim \delta_{\text{DM}}$   
at scales larger than  
the Jeans length  $\sim 1 \text{ com Mpc}$

Viel et al. 2005

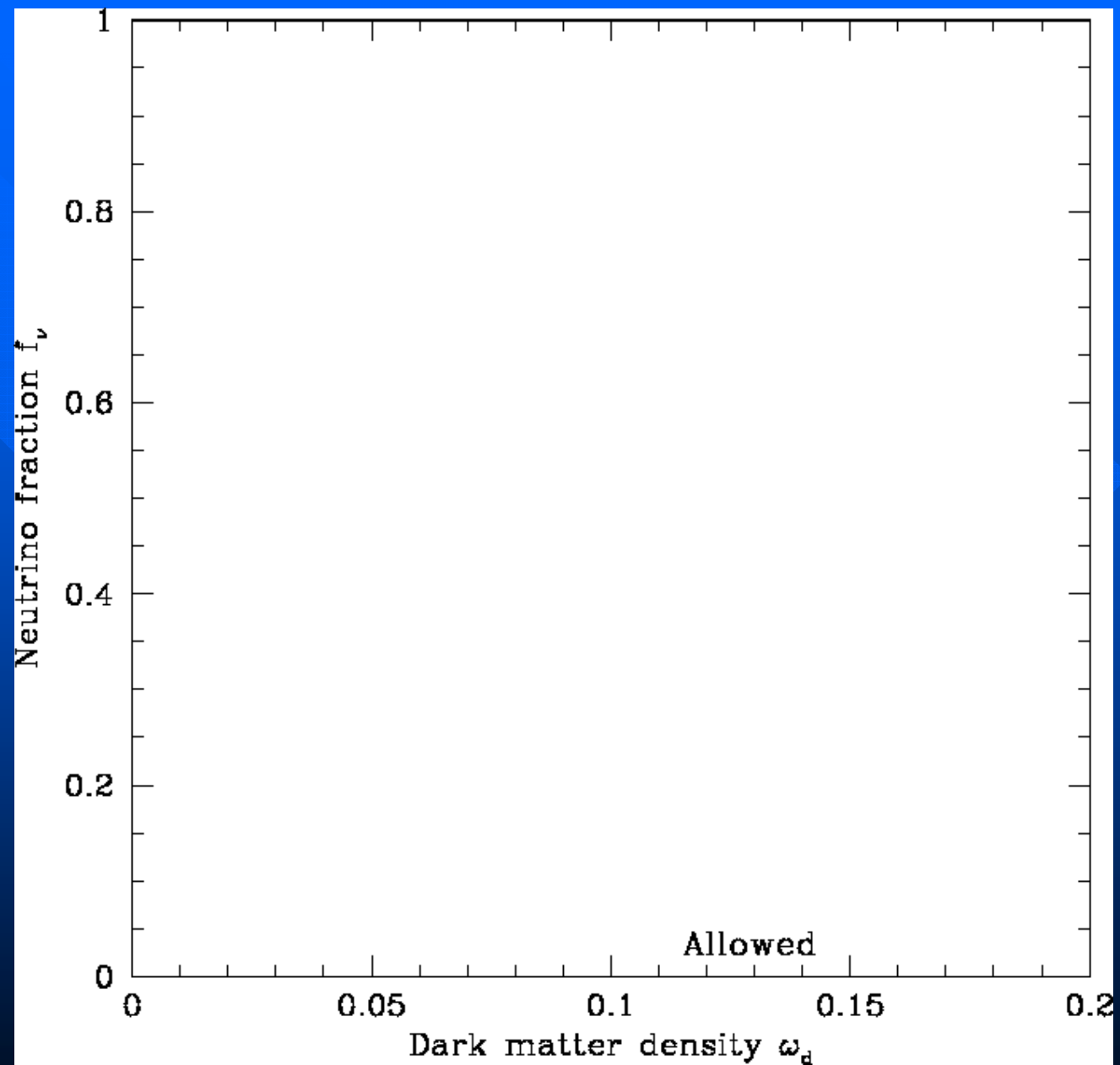


# 3D Lyman-alpha Power Spectrum

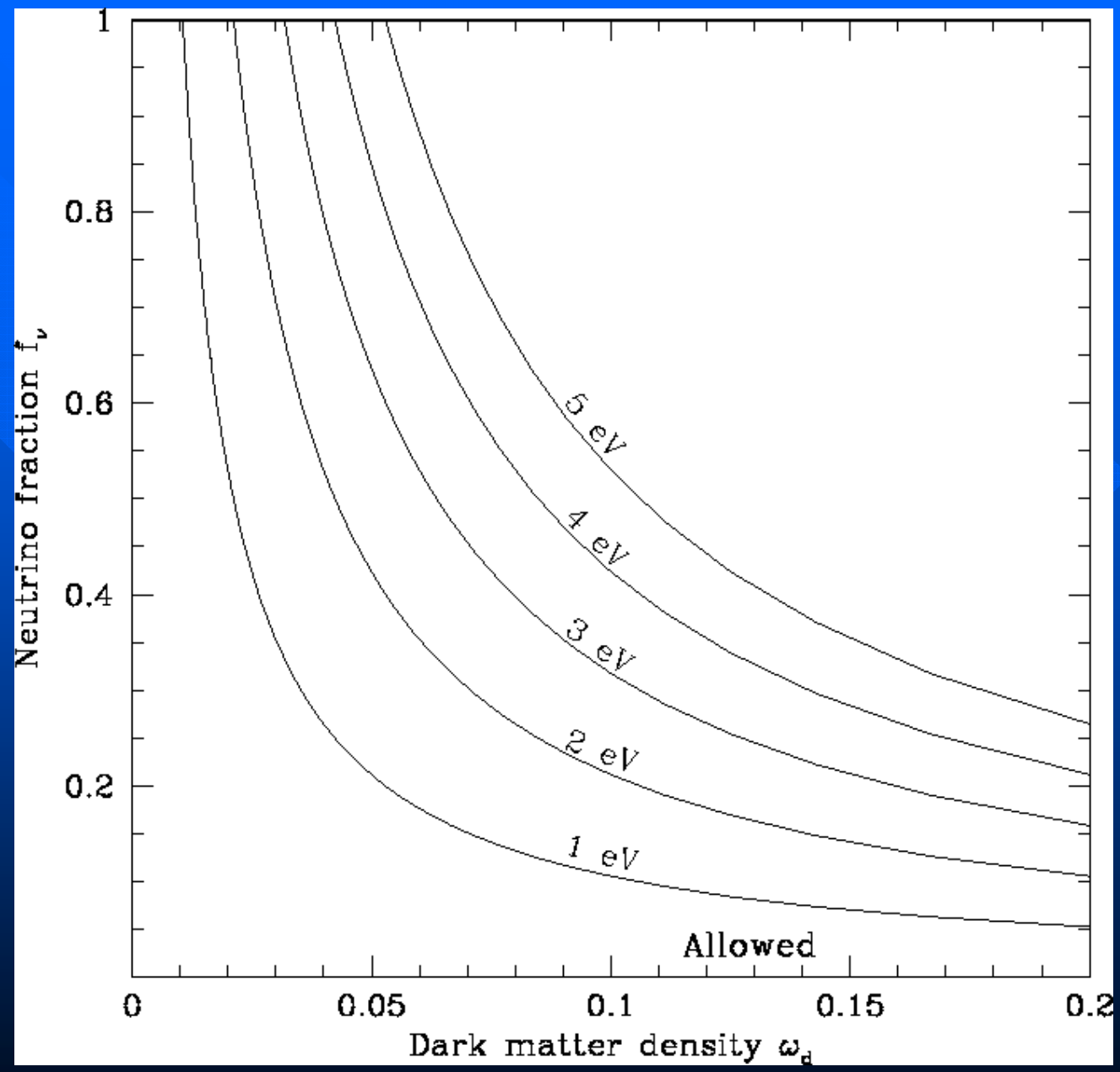
**Very sensitive**  
because at **small**  
**scales**, but quite  
**model-dependent**

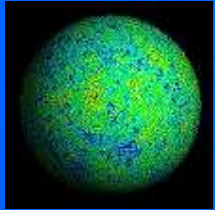


- Run many simulations with CDM-like 3D spectra
- Extract 1D Flux power spectra from each simulation
- Fit amplitude and slope of power at 1 Mpc

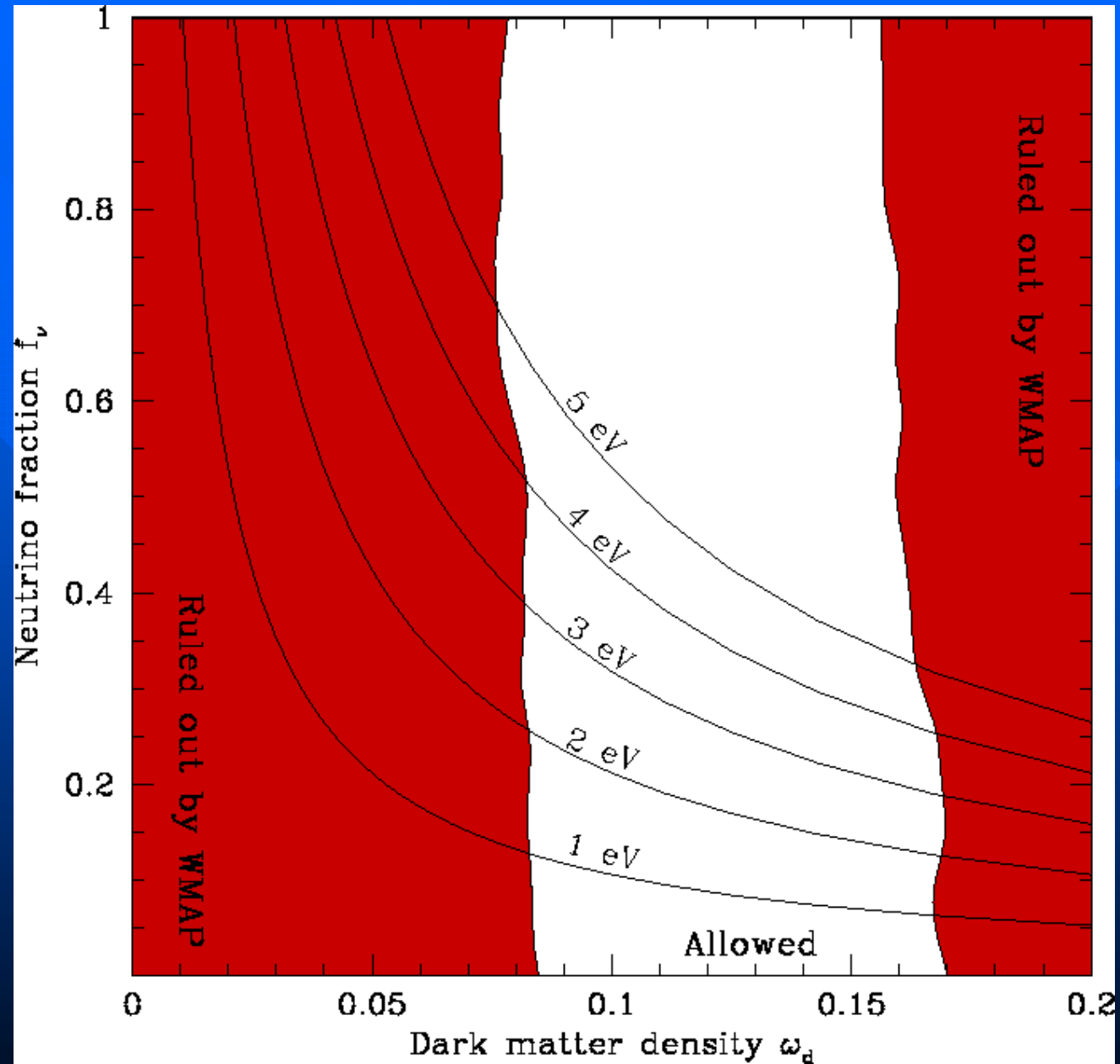


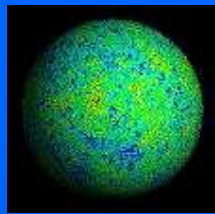






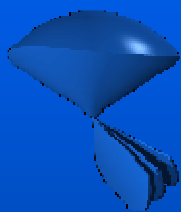
CMB



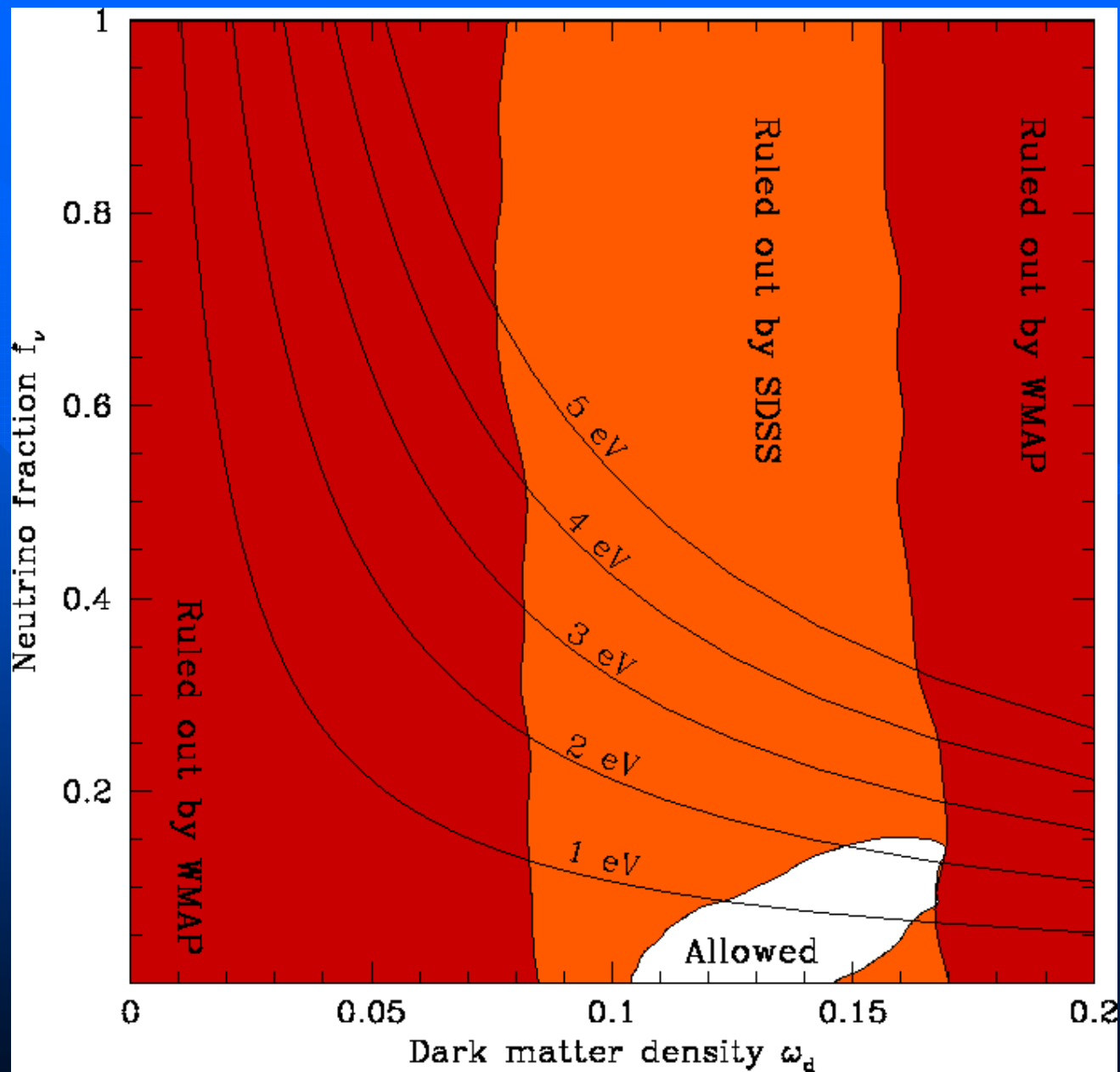


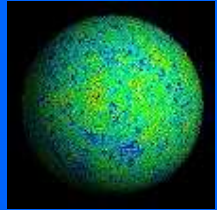
CMB

+

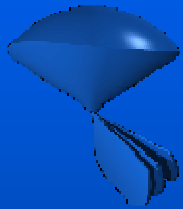


LSS

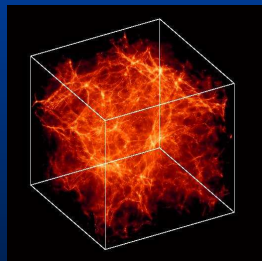




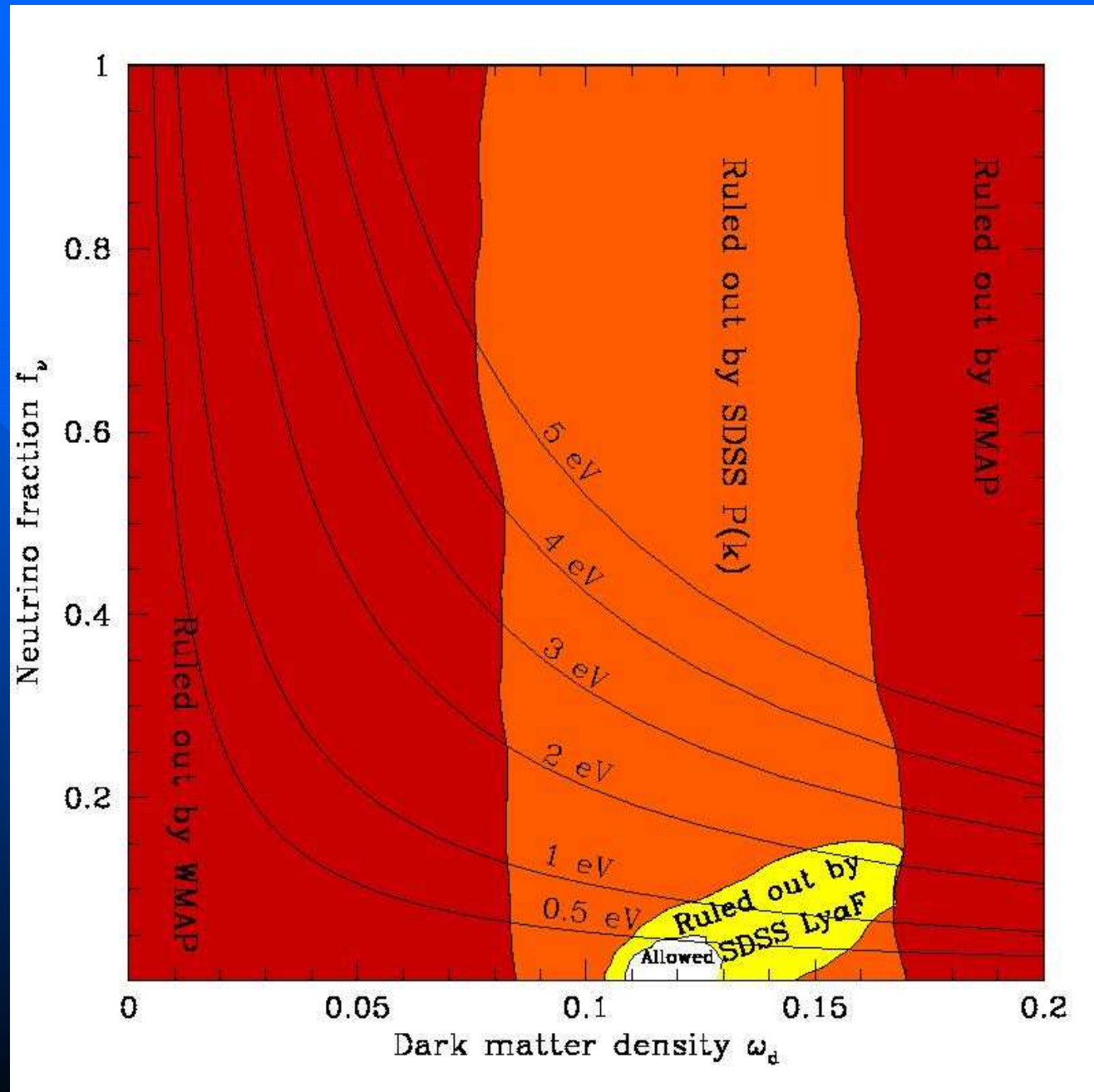
CMB



LSS



Ly $\alpha$ F



# A partial summary of neutrino mass from cosmology

<b>Data</b>	<b>Authors</b>	<b><math>\Sigma m_i</math></b>
2dFGRS	Elgaroy et al 2002	<1.8 eV
WMAP+2dF+..	Spergel et al. 2003	<0.7 eV
WMAP+2dF	Hannestad 2003	<1.0 eV
SDSS+WMAP	Tegmark et al. 2004	<1.7 eV
WMAP+2dF	Crotty et al. 2004	<1.0 eV
WMAP+SDSS Lya	Seljak et al. 2004	<b>&lt;0.43 eV</b>
Clusters+WMAP	Allen et al. 2004	$0.56^{+0.30}_{-0.26}$ eV

**All upper limits 95%, but different assumed priors!**

# Conclusions

- Cosmological constraints on neutrino mass ( $\leq 1\text{eV}$  **total**) arise from *power spectrum* (but attention to priors)
- Wide variety of techniques/experiments needed to eliminate **systematics/degeneracies**
- Physicists must become familiar with: inflation, CMB, LSS, dark energy, ...

# Future Galaxy Cluster Surveys

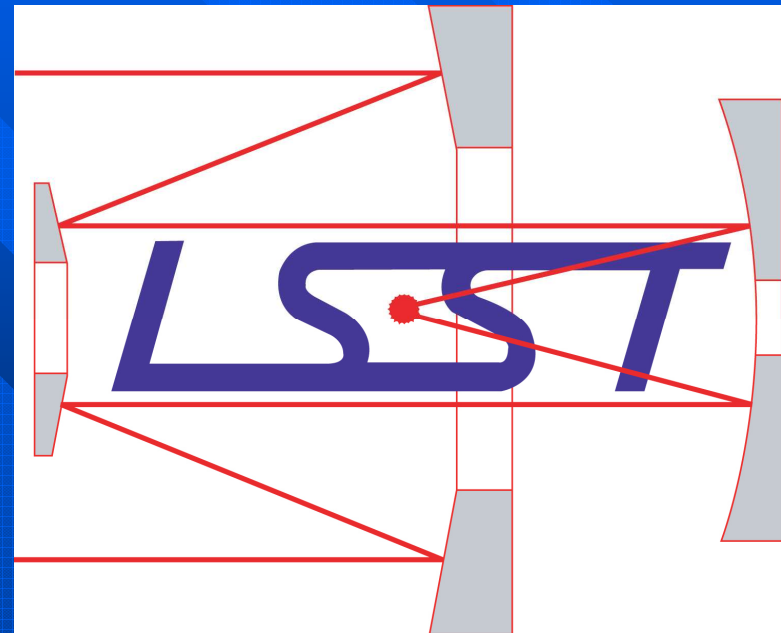
## LSST (Large Synoptic Survey

### Telescope)

A proposed ground-based 8.4-meter telescope detecting galaxy clusters by their weak lensing signals.

Sky coverage: 18000 deg<sup>2</sup>,  
Number of clusters: 200,000;

( $0.1 < z < 1.4$ ,  $M_{\min} = 10^{13.7} h^{-1} M_{\text{sun}}$ )



# Future CMB Surveys

## Planck

Measurement of TT, EE and TE  
in three frequency bands.

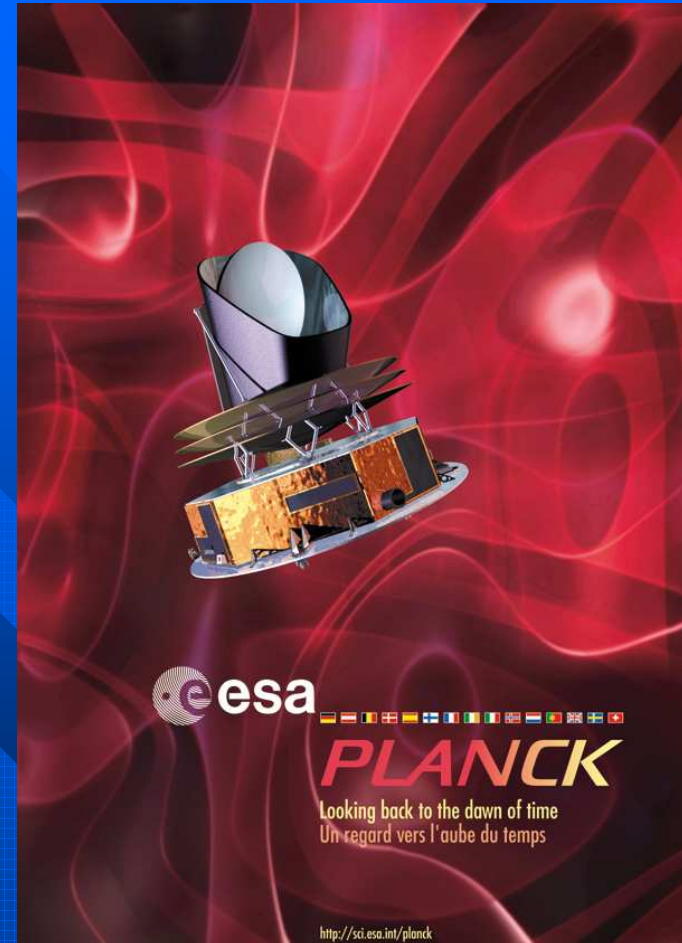
Constraints from CMB  
(unlensed) alone ( $1\sigma$ ):

$$\Delta(w_a) \approx 1.0,$$

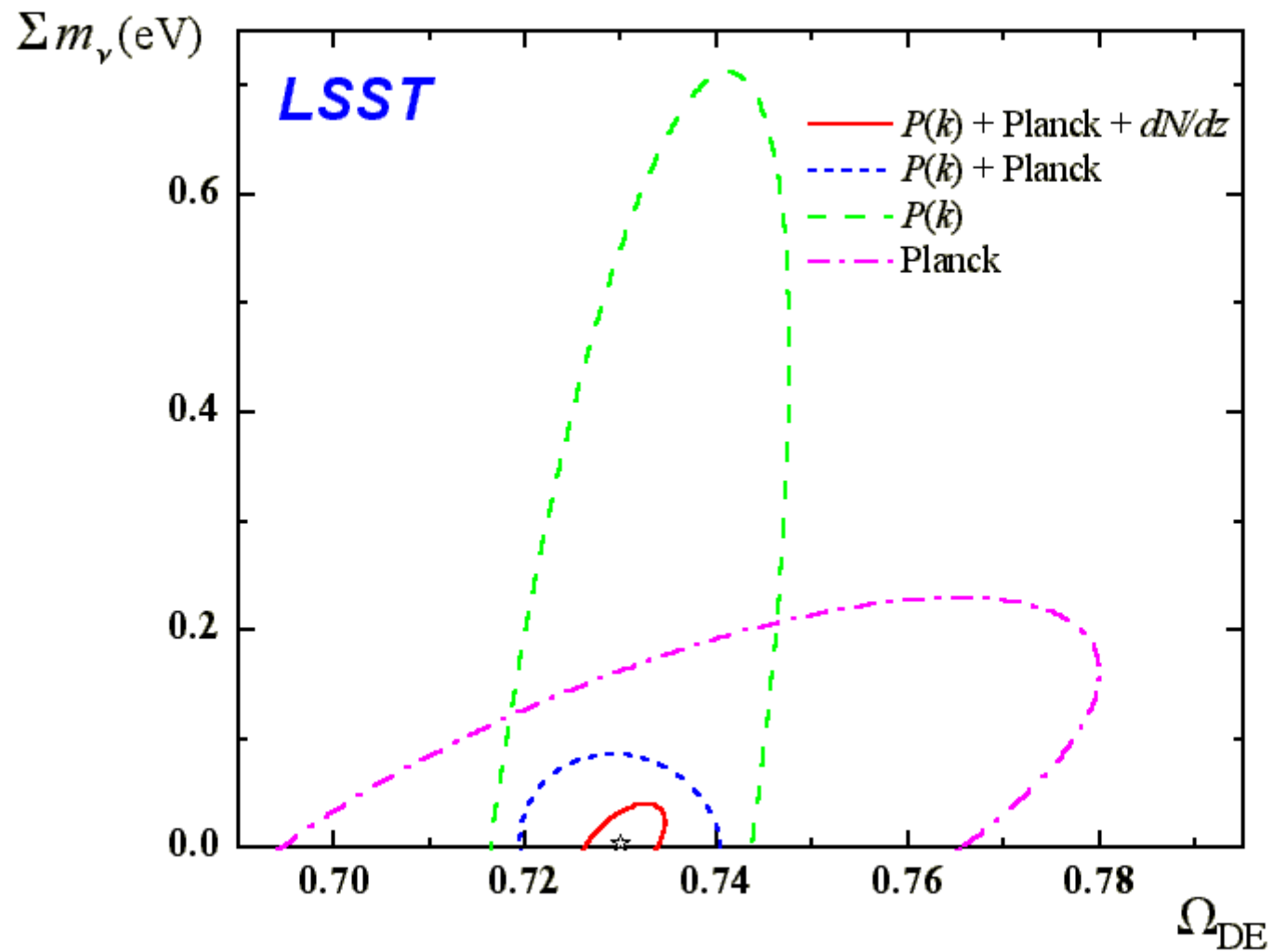
$$\Delta(\Sigma m_\nu) \approx 0.2\text{eV};$$

*e.g.* Eisenstein, Hu & Tegmark  
(1999)

Constraints from clusters will be complementary to those from cosmic microwave background (CMB) anisotropy measurement.

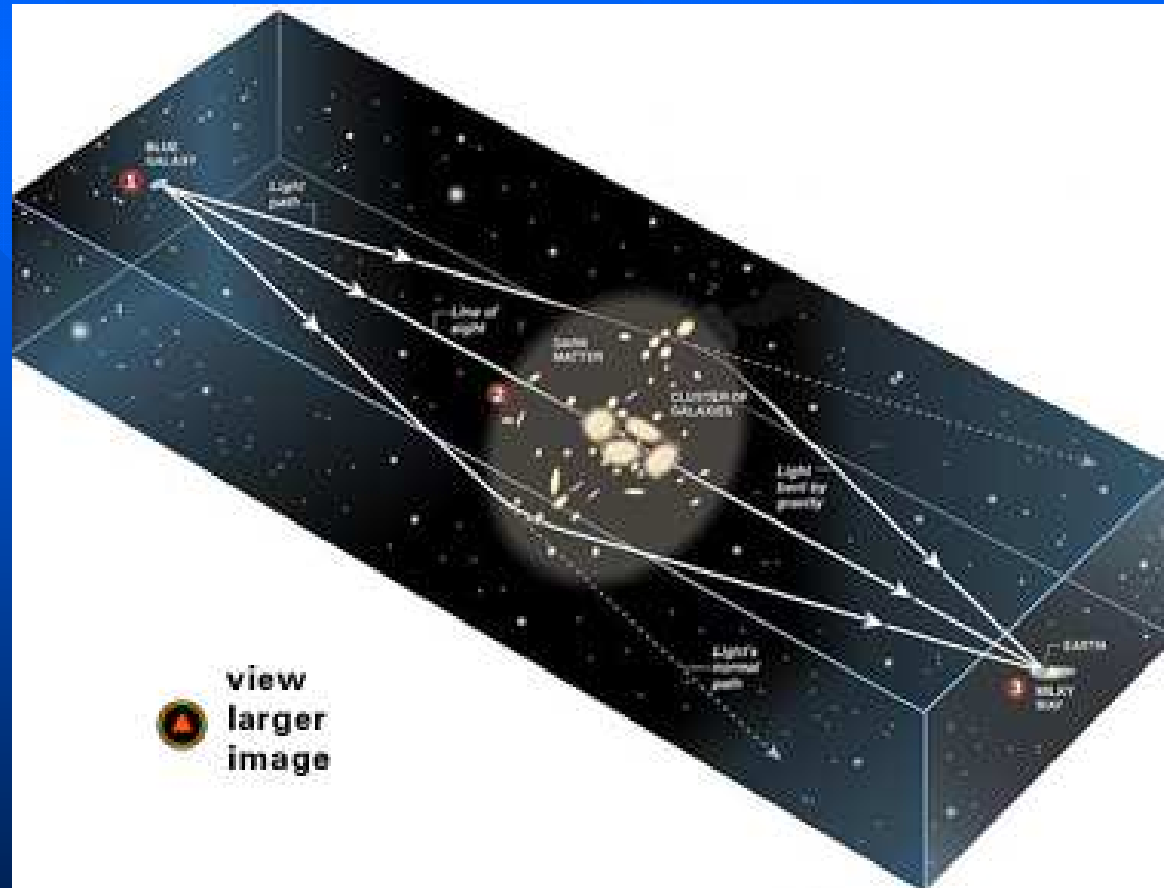






$\Delta(\Sigma m_\nu) \sim 0.04$  eV. (LSST + Planck)

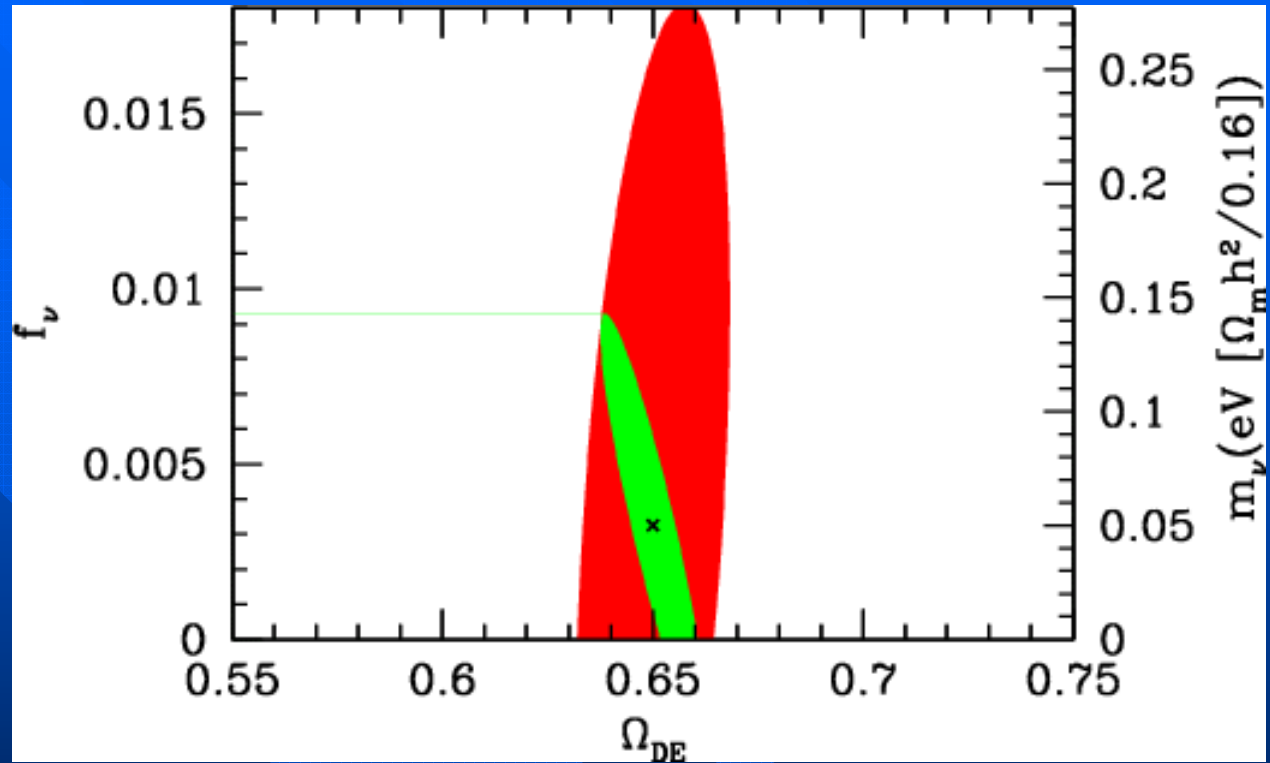
# Weak Gravitational Lensing



**Unlike galaxy surveys and Lyman alpha, lensing directly probes mass distribution!**

# Weak Lensing

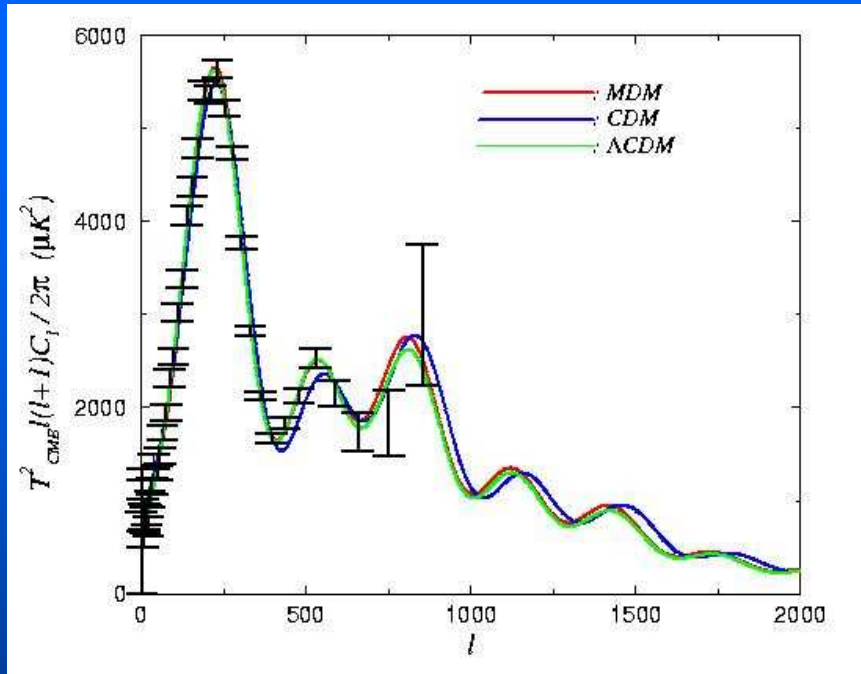
- Measure power spectrum AND/OR measure growth of spectrum at late time
- Sensitive to neutrino mass AND dark energy
- Ergo, accelerator neutrino experiments will teach us about dark energy!



Abazajian & Dodelson (2002)

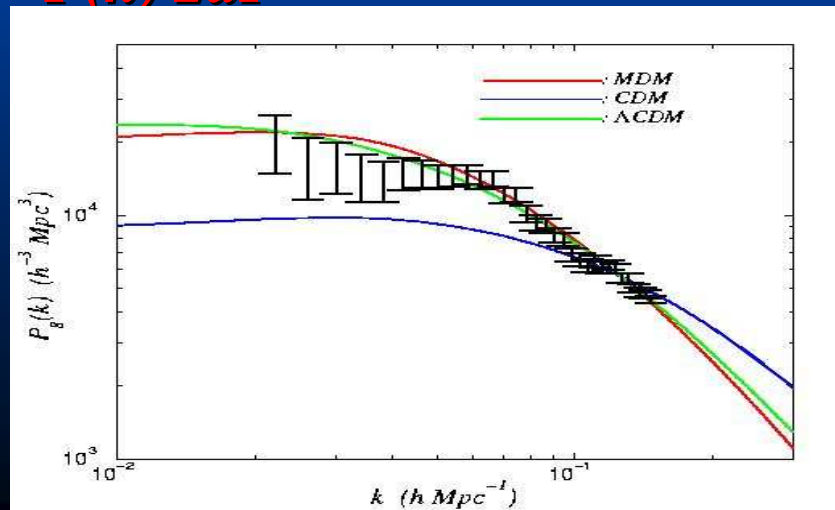
WMAP

# Mixed Dark Matter?



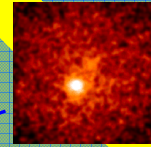
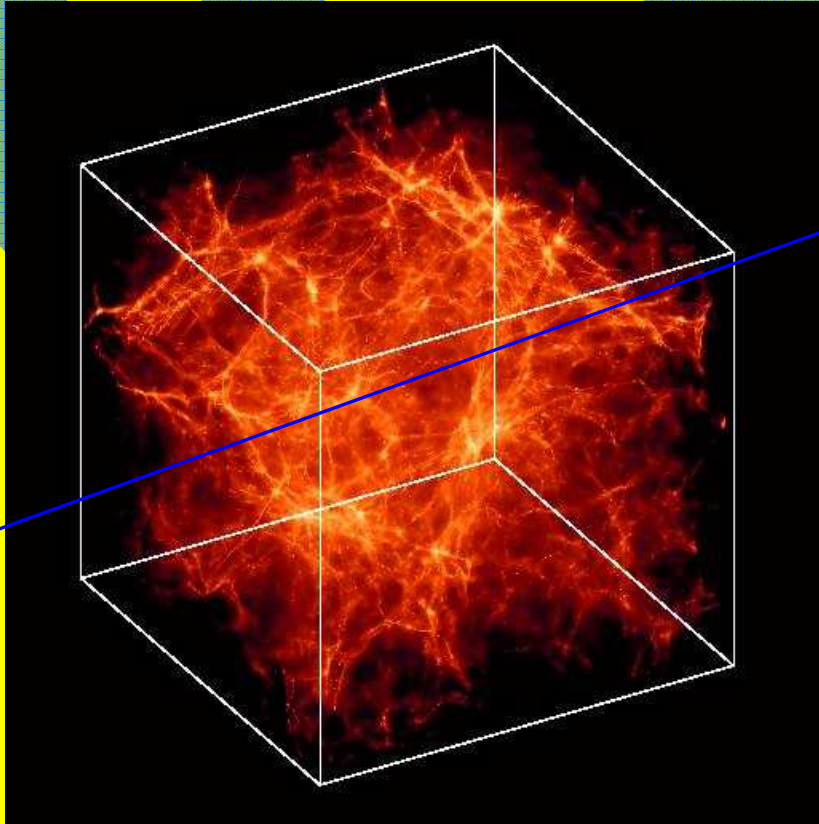
- $\Omega_m=1, \Omega_v=0.2, h=0.45$
- \* Consistent with 2dF.
- \* To fit WMAP, a **break** is required in the Primordial power-spectrum (e.g. Blanchard et al. 2003).

$P(k)$  2dF

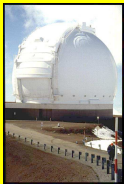


- \* Also at odds with HST's  $H_0$ , SNIa, cluster evolution and baryon fraction.

Elgaroy & Lahav, 2003



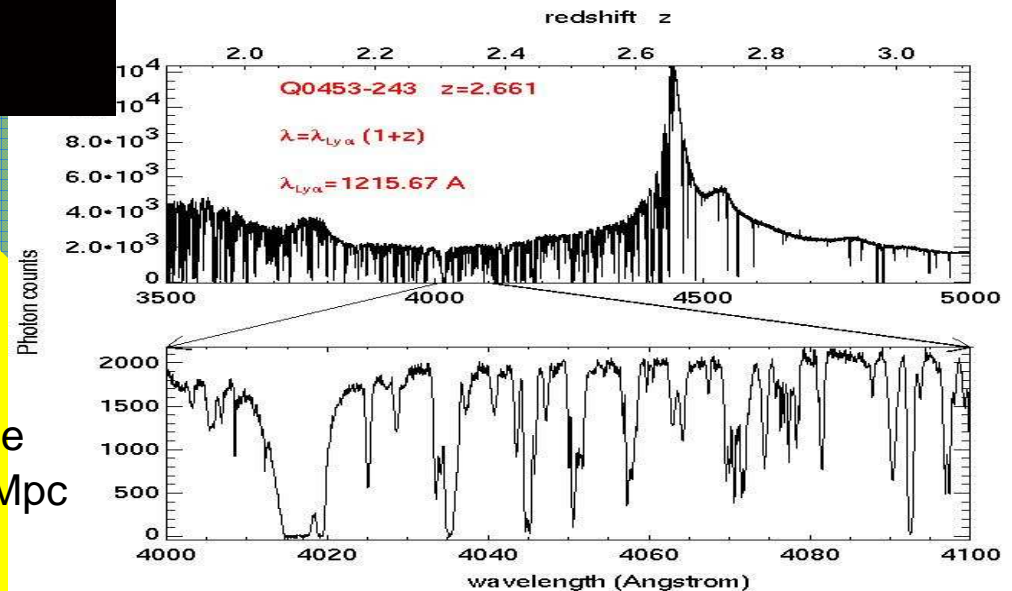
80 % of the baryons at  $z=3$   
are in the Lyman- $\alpha$  forest  
(Rauch 1998)



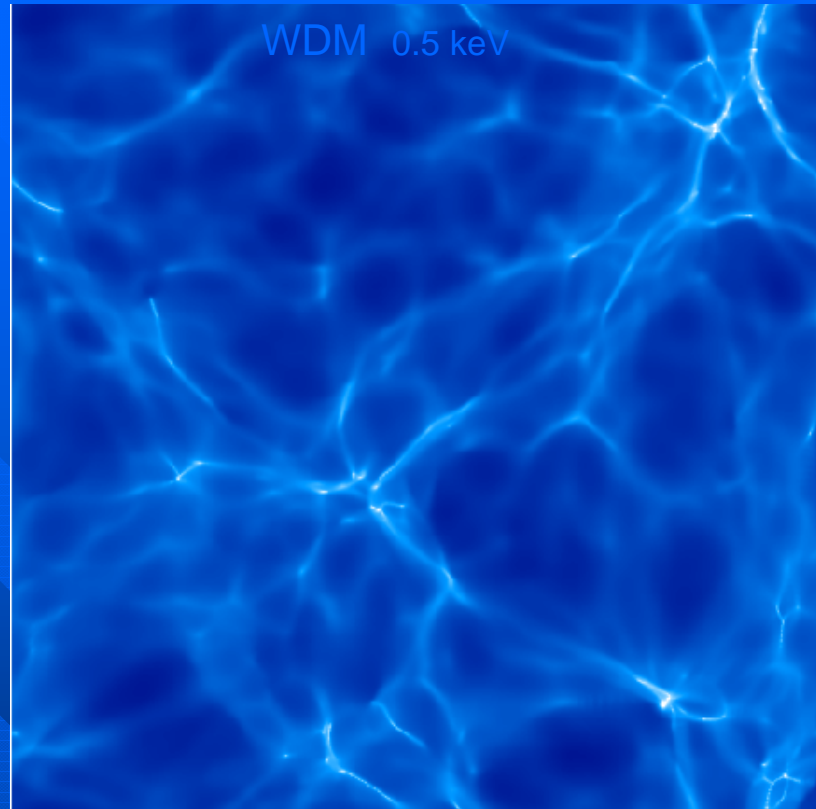
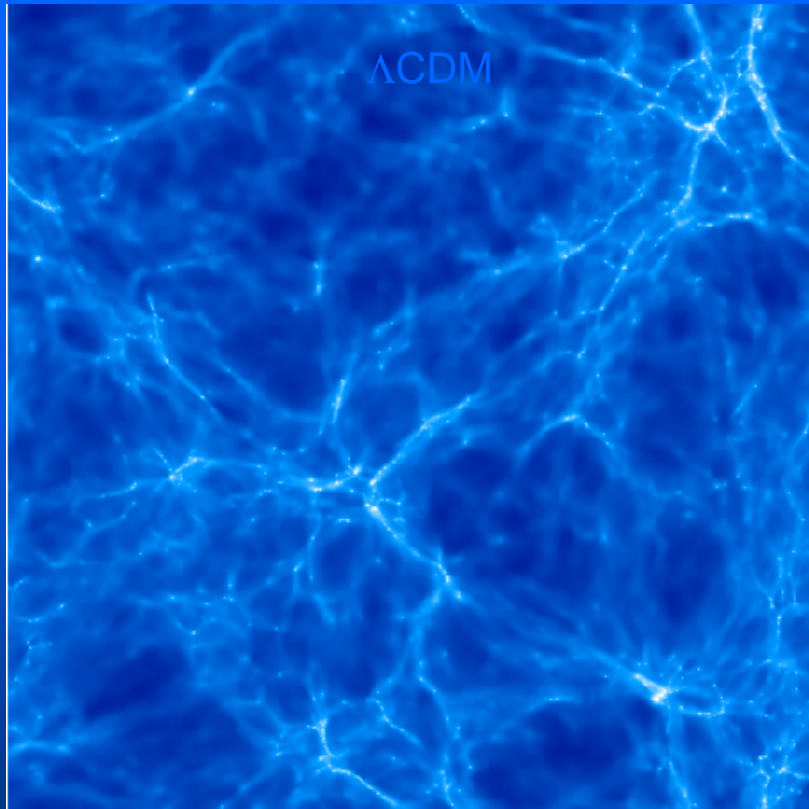
baryons as tracer of the dark  
matter density field

$\delta_{\text{IGM}} \sim \delta_{\text{DM}}$  at scales larger than the  
Jeans length  $\sim 1$  com Mpc

$$\tau \sim (\delta_{\text{IGM}})^{1.6} T^{-0.7} / \Gamma$$



# Cosmological implications: Warm Dark Matter particles-I



30 comoving Mpc/h  $z=3$  gas distribution

In general

$$k_{\text{FS}} \sim 5 \left( T_{\text{v}}/T_{\text{x}} \right) (m \times 1\text{keV}) \text{ Mpc}^{-1}$$

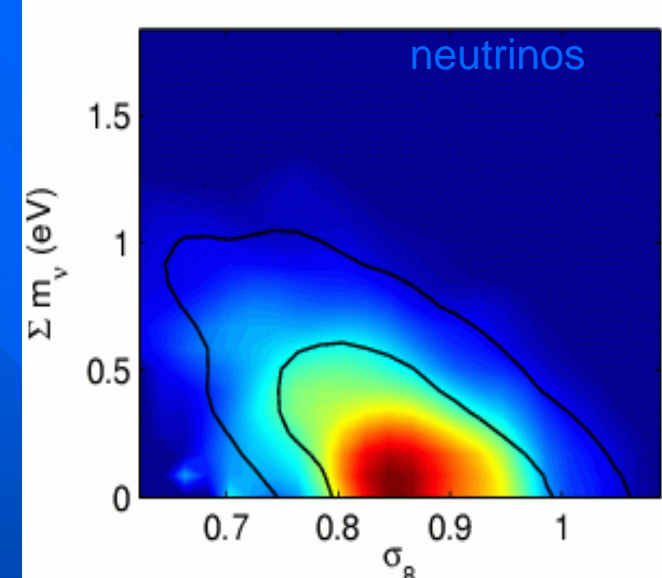
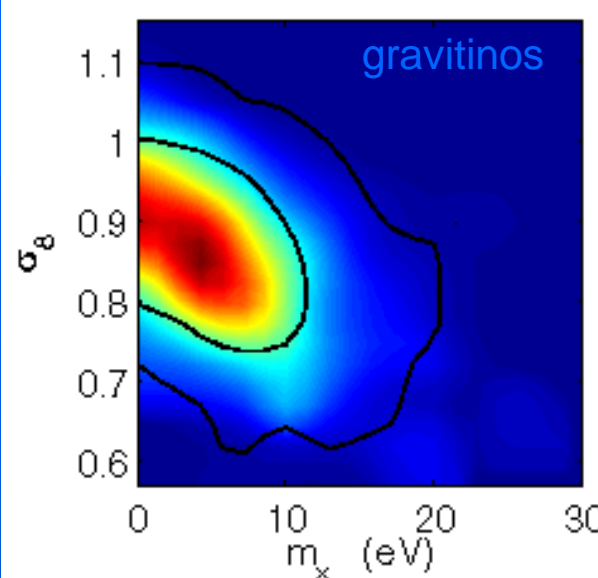
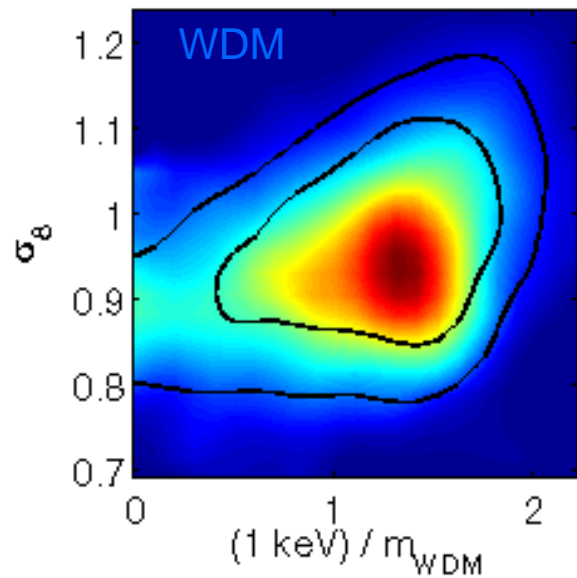


Set by relativistic degrees of freedom at decoupling

if light gravitinos

$$k_{\text{FS}} \sim 1.5 (m \times 100\text{eV}) \text{ h/Mpc}$$

# Cosmological implications: WDM, gravitinos, neutrinos



$m_{\text{WDM}} > 550 \text{ eV}$  ( $2\sigma$ )  
 $> 2\text{keV}$  sterile neutrino ( $2\sigma$ )

$m_{\text{grav}} < 16 \text{ eV}$  ( $2\sigma$ )

$\Sigma m_\nu \text{ (eV)} = 0.33 \pm 0.27$   
 WMAP + 2dF + LY $\alpha$

SDSS gets  $< 0.45$   
 at 95% C.L.  
 (more observables)

	$\Lambda$ WDM	$\Lambda$ CWDM
$\Omega_x h^2$	$0.124 \pm 0.015$	$0.149 \pm 0.019$
$\Omega_B h^2$	$0.024 \pm 0.001$	$0.024 \pm 0.001$
$h$	$0.72 \pm 0.06$	$0.71 \pm 0.06$
$\tau$	$0.18 \pm 0.09$	$0.17 \pm 0.08$
$\sigma_8$	$0.96 \pm 0.08$	$0.86 \pm 0.09$
$n$	$1.01 \pm 0.04$	$1.00 \pm 0.04$
$\alpha \text{ (Mpc}/h)$	$0.06 \pm 0.03$	—
$f_x$	—	$0.05 \pm 0.04$

Set limits on the scale of  
 Supersymmetry breaking  
 if gravitino is the LSP

$\Lambda_{\text{susy}} < 260 \text{ TeV}$

# Neutrinos and the large scale structure

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**Bologna, 17th June 2005**